

2005

Infrared-active phonons of $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($x = 0-0.8$)

Feng Gao

University of Wollongong, fgao@uow.edu.au

Xiaolin Wang

University of Wollongong, xiaolin@uow.edu.au

M Mehdi Farhoudi

University of Wollongong, mmf02@uow.edu.au

R. A. Lewis

University of Wollongong, roger@uow.edu.au

S. X. Dou

University of Wollongong, shi@uow.edu.au

Follow this and additional works at: <https://ro.uow.edu.au/engpapers>



Part of the [Engineering Commons](#)

<https://ro.uow.edu.au/engpapers/43>

Recommended Citation

Gao, Feng; Wang, Xiaolin; Farhoudi, M Mehdi; Lewis, R. A.; and Dou, S. X.: Infrared-active phonons of $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($x = 0-0.8$) 2005.
<https://ro.uow.edu.au/engpapers/43>

INFRARED-ACTIVE PHONONS OF $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($x=0-0.8$)

F. Gao, X. L. Wang, M. Farhoudi, R. A. Lewis and S. X. Dou
Institute for Superconducting & Electronic Materials, University of Wollongong,
Wollongong, NSW2522, Australia

We report an experiment study of infrared active phonons and present a comparative analysis of infrared spectra of polycrystalline $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($x=0-0.8$).

Perovskite containing rare earth and transitional metal cation on A and B sites (with general formula ABO_3) is a class of material which exhibits interesting and useful electronic and magnetic properties. The cobaltite/manganite perovskites $\text{A}(\text{Mn}_{1/2}\text{Co}_{1/2})\text{O}_3$, where A is a lanthanide, exhibit ferromagnetism. The ferromagnetism arises from the superexchange interaction between Co^{2+} and Mn^{4+} ions mediated by oxygen [1]. The Curie temperature, T_C , is relatively high and decreases systematically with decreasing radius of the ion on the A site [2]. The magnetic behavior of rare-earth cobaltite/manganite perovskite compounds is therefore seen to be intimately related to their structure. The magnetic behavior of $\text{A}(\text{Mn}_{1/2}\text{Co}_{1/2})\text{O}_3$ compounds having larger ions (La, Pr, Nd, Sm, Eu) and those having smaller ions (Gd, Tb, Dy, Y, Ho) on the A site has been investigated. Also Phonon modes and phonon energies in cobaltite/manganites, have been determined by far-infrared spectroscopy. The phonon energies systematically shift and split, and new modes appear as the mass of the lanthanide is increased through the series A = La, Nd, Dy, Ho, Yb. This behavior of the phonon modes correlates with the magnetic properties of this series of compounds, in particular with the appearance of metamagnetism for the compounds with smaller ions on the A site [3].

Polycrystalline samples $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ were synthesized by the standard solid-state reaction method. For the samples with $x=0-0.8$ doping level, x-ray diffraction show that structures are orthorhombic and HoMnO_3 is hexagonal. The lattice parameter increases as Co doping level decreases. The ferromagnetism are gradual suppressed when doping level greater and less than $x=0.5$. The polycrystalline samples diluted in CsI have been finely milled and pressed into pellets for infrared transmission measurement. We performed infrared spectroscopy on the $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ samples. For each sample we measured the ratio $I(\omega)/I_0(\omega)$, where $I(\omega)$ is the transmitted intensity through the pellet containing the holmium manganese oxide and that $I_0(\omega)$ is the transmitted intensity through pure CsI pellet. The infrared transmission spectra of polycrystalline $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ at room temperature are shown in Figure 1. The normalized spectra of optical density is proportional to the optical conductivity $\sigma(\omega)$ [4,5]

$$O_d(\omega) = \ln[I_0(\omega)/I(\omega)] \propto \sigma(\omega). \quad (1)$$

The energies of phonons $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ samples can be obtained from spectra of optical density. The spectra exhibit numerous peaks associated with the infrared active phonons modes. To extract information on phonons we fit the spectra of optical density using a sum of

noninteracting harmonic oscillators (lorentz model)[6,7]. The minimum number of oscillators obtaining a reliable fit for samples $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ will be reported.

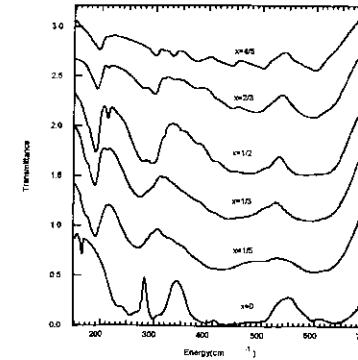


Fig.1. Transmission spectra of polycrystalline $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($X=0-0.8$) at room temperature.

The optical modes of the ideal cubic structure have the irreducible representation $\Gamma = 3F_{1u}$ (infrared active) + F_{2u} (infrared inactive). Last [8] identifies the bands in order of increasing energy as external A-(BO₃) vibration, O-B-O bending, and B-O stretching. As the crystal symmetry is reduced, more optical modes are expected. There are 25 infrared active modes for orthorhombic (D_{2h}^{16} , Pnma) symmetry [9]. Consequently, the four broad bands for $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ ($x=0.2-0.8$) at $\omega_1 \sim 190 \text{ cm}^{-1}$, $\omega_2 \sim 270 \text{ cm}^{-1}$, $\omega_3 \sim 420 \text{ cm}^{-1}$, $\omega_4 \sim 600 \text{ cm}^{-1}$, are assigned to the external, torsional, bending and stretching modes, respectively. The behavior of four main bands of $\text{HoMn}_{1-x}\text{Co}_x\text{O}_3$ with different Co element doping level will be analysed in detail.

References

- 1 X. L. Wang et al., *Supercond. Sci. Technol.*, **15**, 427(2002).
- 2 I. O. Troyanchuk et al., *Mater. Res. Bull.* **32**, 67(1997).
- 3 F. Gao et al., *J. solid State Chemistry* **160**, 350(2001).
- 4 A. Paolone et al., *Physica (Amsterdam)*, **244B**, 33(1998).
- 5 P. Calvani et al., *Phys. Rev. Lett.* **81**, 4504(1998).
- 6 I. Fedorov et al., *Phys. Rev. B*, **60**, 11875(1999).
- 7 G. Burns, *Solid State Physics*(Academic Press, New York, 1985).
- 8 J. T. Last, *Phys. Rev.* **105**, 1740(1957).
- 9 Q. Williams et al., *J. Geophys. Res.*, **27**, 8116(1987).