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### Trap-assisted transport and non-uniform charge distribution in sulphur-rich PbS colloidal quantum dot-based solar cells with selective contacts

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# Trap-assisted transport and non-uniform charge distribution in sulphur-rich PbS colloidal quantum dot-based solar cells with selective contacts

## Abstract

This study reports evidences of dispersive transport in planar PbS colloidal quantum dots heterojunction-based devices as well as the effect of incorporating a MoO<sub>3</sub> hole selective layer on the charge extraction behavior. Steady state and transient characterization techniques are employed to determine the complex recombination processes involved in such devices. The addition of a selective contact drastically improves the device efficiency up to 3.15 % (especially through the photocurrent and series resistance) and extends the overall charge lifetime by suppressing the main first-order recombination pathway observed in device without MoO<sub>3</sub>. The lifetime and mobility calculated for our sulphur-rich PbS-based devices are similar to previously reported values in lead-rich quantum dots-based solar cells. Nevertheless, strong Shockley-Read-Hall mechanisms appears to keep restricting charge transport, the equilibrium voltage taking more than 1 ms to be established.

## Keywords

contacts, cells, solar, dot, quantum, colloidal, pbs, rich, sulphur, distribution, selective, charge, trap, uniform, non, transport, assisted

## Disciplines

Engineering | Physical Sciences and Mathematics

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## SUPPORTING INFORMATION

# Trap-Assisted Transport and Non-Uniform Charge Distribution in Sulphur-Rich PbS Colloidal Quantum Dot-based Solar Cells with Selective Contacts

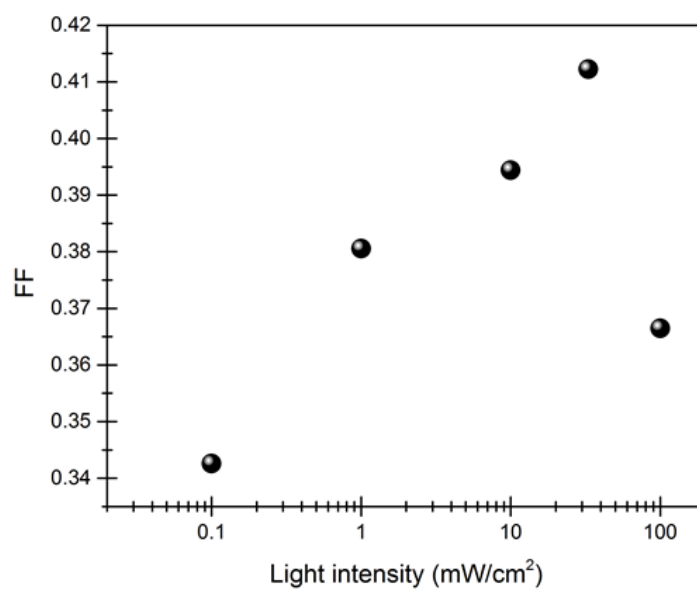
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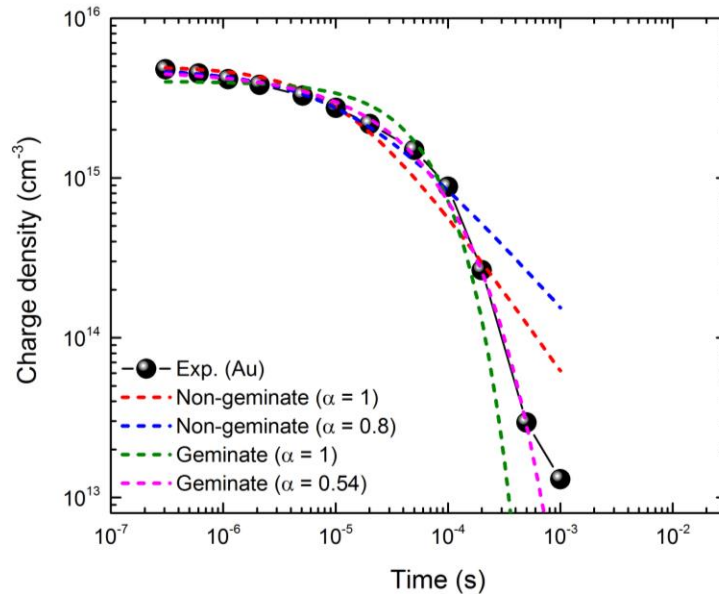
**Figure S1.** Fill factor (FF) as a function of the light intensity.

## Fitting of dispersive exponential

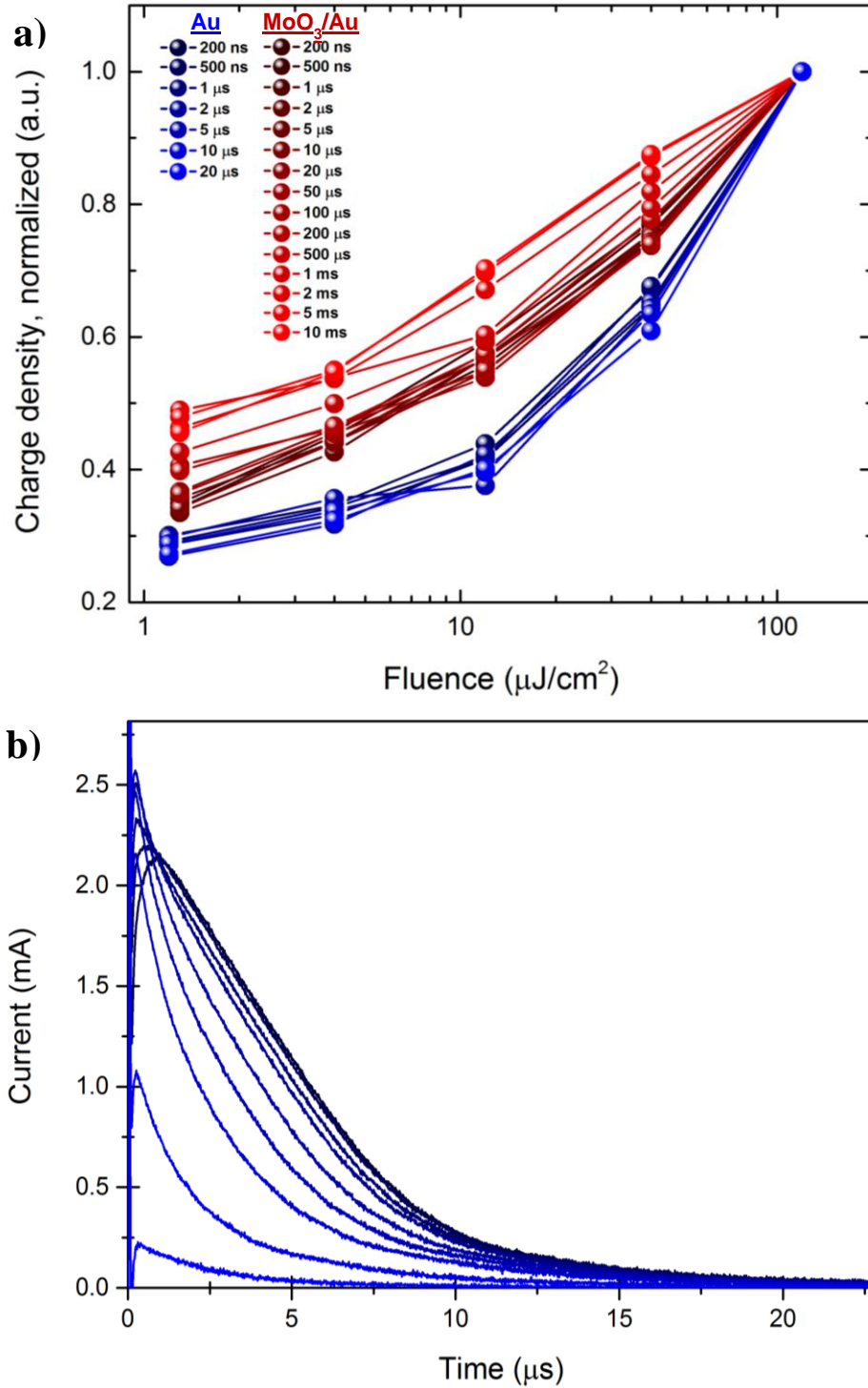
The fitting functions are:

- Geminate:  $n(t) = n_0 e^{-\left(\frac{t}{\tau}\right)^\alpha}$
- Non-geminate:  $n(t) = \frac{n_0}{1+n_0(\beta t)^\alpha}$

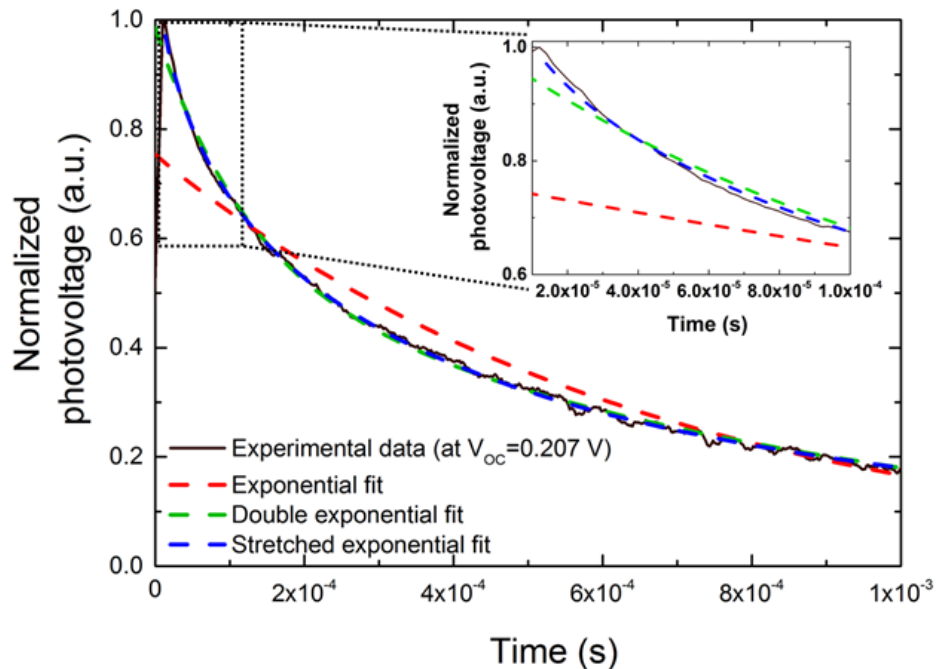
where  $n_0$  is the initial concentration of free photogenerated charges and  $\tau$ ,  $\beta$ , and  $\alpha$  are the geminate recombination characteristic time, the non-geminate recombination coefficient and the dispersive factor. These two functions are obtained from solving the continuity reaction:  $\frac{dn}{dt} = \frac{dp}{dt} = -\beta n(t)p(t)$ . In the case of geminate recombination, the equation can be simplified as the density of one charge carrier can be considered constant (e.g.,  $n(t) = n$  and  $\beta n = \frac{1}{\tau}$ ), thus leading to a solution following a monoexponential decay. Typical non-geminate, assuming charge quasi-neutrality ( $n(t) \approx p(t)$ ), results in a hyperbolic solution. The dispersive factor arises from taking into account the dispersive media where mobile charges have a high probability of encountering deep traps with a slow release time which compete with other characteristic transport times (diffusion, injection/transfer). For a non-dispersive random walk, the number of sites visited by the mobile charges is linear ( $S(t) \propto t$ ). On the other, considering a dispersion in the spatial separation between the localized states and a dispersion in the potential barriers between these sites, Scher and Montroll demonstrated that the waiting time distribution function could be well-represented by a power law  $S(t) \propto t^\alpha$  where  $\alpha = T/T_0$  and  $T_0$  is the characteristic temperature of the distribution.<sup>1,2</sup> In the limit where  $T_0 = T$  (or  $\alpha = 1$ ), the distribution is inexistent and the solution is non-dispersive.



**Figure S2.** Experimental time-resolved charge density decay of the pixel without selective contact along with various fitting functions.



**Figure S3.** a) Comparison between normalized charge density in function of the light intensity, highlighting the fact that the decays are similar up to 20  $\mu\text{s}$ . b) Current extraction curves for the device without selective contact under  $40 \mu\text{J}/\text{cm}^2$  at various delay time (from 200 ns to 100  $\mu\text{s}$ ) highlighting the linear nature of the decay at fast delays.



**Figure S4.** Typical TPV signal fitted with various exponential functions. Inset: high resolution of the early stage of the decay, highlighting the difference between the double exponential and the stretched exponential.

## References

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S2. Nelson J.; Chandler R. E. Random walk models of charge transfer and transport in dye sensitized systems. *Coord. Chem. Rev.* 2004, 248, 1181-1194.