

University of Wollongong

Research Online

Faculty of Informatics - Papers (Archive)

Faculty of Engineering and Information
Sciences

1-1-2012

Surface diffusion flow near spheres

Glen Wheeler

University of Wollongong, glenw@uow.edu.au

Follow this and additional works at: <https://ro.uow.edu.au/infopapers>



Part of the [Physical Sciences and Mathematics Commons](#)

Recommended Citation

Wheeler, Glen: Surface diffusion flow near spheres 2012, 131-151.
<https://ro.uow.edu.au/infopapers/1878>

Research Online is the open access institutional repository for the University of Wollongong. For further information contact the UOW Library: research-pubs@uow.edu.au

Surface diffusion flow near spheres

Abstract

We consider closed immersed hypersurfaces evolving by surface diffusion flow, and perform an analysis based on local and global integral estimates. First we show that a properly immersed stationary ($\Delta H \equiv 0$) hypersurface in \mathbb{R}^3 or \mathbb{R}^4 with restricted growth of the curvature at infinity and small total tracefree curvature must be an embedded union of umbilic hypersurfaces. Then we prove for surfaces that if the L^2 norm of the tracefree curvature is globally initially small it is monotonic nonincreasing along the flow. We also derive pointwise estimates for all derivatives of the curvature assuming that its L^2 norm is locally small. Using these results we show that if a singularity develops the curvature must concentrate in a definite manner, and prove that a blowup under suitable conditions converges to a nonumbilic embedded stationary surface. We obtain our main result as a consequence: the surface diffusion flow of a surface initially close to a sphere in L^2 is a family of embeddings, exists for all time, and exponentially converges to a round sphere.

Keywords

surface, near, diffusion, spheres, flow, ERA2015

Disciplines

Physical Sciences and Mathematics

Publication Details

Wheeler, G. (2012). Surface diffusion flow near spheres. *Calculus of Variations and Partial Differential Equations*, 44 (1-2), 131-151.

SURFACE DIFFUSION FLOW NEAR SPHERES

GLEN WHEELER

ABSTRACT. We consider closed immersed hypersurfaces evolving by surface diffusion flow, and perform an analysis based on local and global integral estimates. First we show that a properly immersed stationary ($\Delta H \equiv 0$) hypersurface in \mathbb{R}^3 or \mathbb{R}^4 with restricted growth of the curvature at infinity and small total tracefree curvature must be an embedded union of umbilic hypersurfaces. Then we prove for surfaces that if the L^2 norm of the tracefree curvature is globally initially small it is monotonic nonincreasing along the flow. We also derive pointwise estimates for all derivatives of the curvature assuming that its L^2 norm is locally small. Using these results we show that if a singularity develops the curvature must concentrate in a definite manner, and prove that a blowup under suitable conditions converges to a nonumbilic embedded stationary surface. We obtain our main result as a consequence: the surface diffusion flow of a surface initially close to a sphere in L^2 is a family of embeddings, exists for all time, and exponentially converges to a round sphere.

1. INTRODUCTION

Let $f : M^n \times [0, T] \rightarrow \mathbb{R}^{n+1}$ be a family of compact immersed hypersurfaces $f(\cdot, t) = f_t : M \rightarrow f_t(M) = M_t$ with associated Laplace-Beltrami operator Δ , unit normal vector field ν , and mean curvature function H . The surface diffusion flow

$$(1) \quad \frac{\partial}{\partial t} f = (\Delta H)\nu,$$

is the chief object of interest for this paper. Equation (1) is a fourth order degenerate quasilinear parabolic evolution problem, the local existence of which is now standard in the literature; see [8] for example. Our goal in this paper is to establish regularity and stability results analagous to the pioneering work of Kuwert & Schätzle [14, 15].

The hallmark geometric characteristics of surface diffusion flow are the following: using $\text{Vol } M_t$ to denote the volume enclosed by M_t in \mathbb{R}^{n+1} we compute

$$(2) \quad \frac{d}{dt} \text{Vol } M_t = \int_M \Delta H d\mu = 0, \text{ and}$$

$$(3) \quad \frac{d}{dt} \int_M d\mu = \int_M H \Delta H d\mu = - \int_M |\nabla H|^2 d\mu \leq 0;$$

so that a manifold evolving by (1) will exhibit conservation of enclosed volume and monotonic decreasing surface area. Further, surface area is preserved exactly when the mean curvature of M_t is constant. This is similar to mean curvature flow (with normal velocity $\partial_t^\perp f = -H$) where surface area is monotonically decreasing, and stationary if M_t is a minimal surface. As mean curvature flow may be used as a tool in studying minimal surfaces, surface diffusion flow may be used to study surfaces of

1991 *Mathematics Subject Classification.* 53C44 and 58J35.

Key words and phrases. global differential geometry and fourth order and geometric analysis and surface diffusion flow and higher order partial differential equations.

constant mean curvature with a prescribed volume. These properties make surface diffusion flow one of the most natural fourth order flows one can consider (with the other candidate being Willmore flow, where $\partial_t^\perp f = \Delta H + H|A^\circ|^2$), and a model problem to be well studied before moving on to more general evolution equations.

Since one may derive (1) by considering the H^{-1} -gradient flow for the area functional (see Fife [9]), one may suspect that a program of study similar to that presented by Kuwert & Schätzle [14, 15] on the gradient flow for the Willmore functional is not appropriate. Indeed, the natural reaction is to guess that surface diffusion flow has more in common with mean curvature flow than Willmore flow. With this in mind, an immediate goal is to establish a lower bound on the surface area of the evolving hypersurface. This would follow from the computation (2) and the isoperimetric inequality, so long as we can show that if f_0 is an embedding, then f_t is also an embedding. The crucial point is to prove (Proposition 12) that under assumption (4) below the curvature is monotonically nonincreasing along the flow. Combining this with a result from Li & Yau [16] finally allows us to conclude that embeddedness is preserved for surface diffusion flows satisfying (4), and that the surface area is uniformly bounded away from zero.

With our analysis in this paper we affirm that the general technique of localised integral estimates and blowup analysis from [14, 15] is effective for surface diffusion flow. Our main result is the following theorem, analagous to Theorem 5.1 in [14].

Theorem 2. *There exists an absolute constant $\epsilon_0 > 0$ such that if $f : M^2 \times [0, T) \rightarrow \mathbb{R}^3$ is a surface diffusion flow with*

$$(4) \quad \int_M |A^\circ|^2 d\mu \Big|_{t=0} \leq \epsilon_0 < 8\pi$$

then $T = \infty$, $f : M^2 \times [0, \infty) \rightarrow \mathbb{R}^3$ is a family of embeddings, and $M_t \rightarrow S^2$ exponentially, where S^2 is a round sphere.

Our larger aim is to instigate a systematic study of the asymptotic behaviour of the surface diffusion flow. One way to view the condition (4) is that the deviation of f from being round is small in an averaged sense. Our result here can then be viewed as a kind of stability of spheres theorem in the L^2 norm. It then becomes natural to query on the global behaviour of f when this deviation is greater than ϵ_0 , or even if it is initially very large. At this time, we do not know of any analytic example which rules out dropping condition (4), or allows us to provide an upper bound for ϵ_0 . As a first step, it would be interesting to determine whether an immersed sphere with image a symmetric figure 8, having zero enclosed volume, shrinks to a point and vanishes in finite time.

The surface diffusion flow has been considered for some time in the literature. First proposed by Mullins [20] in 1957 (two years before he proposed the mean curvature flow), it was originally designed to model the formation of thermal grooves in phase interfaces where the contribution due to evaporation-condensation was insignificant. Some time later, Davi, Gurtin, Cahn and Taylor [4, 6] proposed many other physical models which give rise to the surface diffusion flow. These all exhibit a reduction of free surface energy and conservation of volume; essential characteristics of surface diffusion flow. There are also other motivations for the study of (1). For example, two years later Cahn, Elliot and Novick-Cohen [3] proved that

(1) is the singular limit of the Cahn-Hilliard equation with a concentration dependent mobility. Among other applications, this arises in the modeling of isothermal separation of compound materials.

Analysis of the surface diffusion flow began slowly. Baras, Duchon and Robert [1] showed the global existence of weak solutions for two dimensional strip-like domains. Later, Elliot and Garcke [7] analysed the surface diffusion flow of curves, and obtained local existence and regularity for C^4 -initial curves, and global existence for small perturbations of circles. Escher, Mayer and Simonett [8] gave several numerical schemes for modeling (1), and have also given the only two known numerical examples [18] of the development of a singularity: a tubular spiral and thin-necked dumbbell. Ito [13] showed that convexity will not be preserved under surface diffusion flow, even for smooth, rotationally symmetric, closed, strictly convex initial hypersurfaces. Quite recently, Blatt [2] generalised this and showed loss of convexity and loss of embeddedness for a large class of flows.

There have been many important works on fourth order flows of a slightly different character, from Willmore flow of surfaces to Calabi flow, a fourth order flow of metrics. Significant contributions to the analysis of these flows by the authors Kuwert, Schätzle, Polden, Huisken, Mantegazza and Chruściel [5, 14, 15, 17, 21] are particularly relevant, as the methods employed there are similar to ours here.

Simonett [23] used centre manifold techniques to show that the statement of Theorem 2 holds under the stronger assumption that f_0 is $C^{2,\alpha}$ -close to a round sphere. Our analysis here is completely different, drawing inspiration instead from the work of Kuwert & Schätzle [14, 15] on the Willmore flow of surfaces. The most transparent difference between surface diffusion and Willmore flow is that one lacks the explicit structure of an L^2 gradient flow. While we are able to show that under assumption (4) we have

$$\frac{d}{dt} \int_M |A^\circ|^2 d\mu \leq -\frac{1}{4} \int_M |\Delta H|^2 d\mu,$$

it is completely unknown whether or not this is true for initial data violating (4).

This paper is organised as follows. In Section 2 we derive integral estimates in the case where the tracefree curvature is locally small in L^n , where $n \in \{2, 3\}$. Using these we conclude our first results, curvature estimates and a gap lemma on the stationary solutions of (1) satisfying a small tracefree curvature assumption and restricted growth of the second fundamental form in L^2 at infinity. The key ingredient allowing the classification result to go through in three intrinsic dimensions is a new multiplicative Sobolev inequality. We also outline the proof of a lifespan theorem and interior estimates, which are analogous to Theorem 1.2 in [15] and Theorem 3.5 in [14]. Section 3 contains the blowup construction we will employ, and shows that for surface diffusion flows satisfying (4) the blowup at a finite time curvature singularity is a stationary, nonumbilic surface with $\Delta H \equiv 0$. In Section 4 we show that this is in contradiction with the gap lemma, and thus obtain long time existence for surface diffusion flows satisfying (4). In Section 4 we also examine the global behaviour of such a flow and show exponential convergence to spheres.

We have developed the exposition to be particularly relevant to the case of globally constrained surface diffusion and Willmore flows, where the immersion evolves by

$$(5) \quad \frac{\partial}{\partial t} f = (\Delta H + h)\nu, \quad \text{and} \quad \frac{\partial}{\partial t} f = (\Delta H + H|A^\circ|^2 + h)\nu$$

respectively, for some constraint $h : [0, T) \rightarrow \mathbb{R}$. The flow is non-local, in the sense that the motion of points on M may depend upon global properties of M , such as total curvature, surface area, mixed volumes, or other quantities. Choices of h are motivated with some geometric consideration in mind, and give rise to mixed volume preserving surface diffusion flow, and surface area, volume preserving Willmore flows for example. With suitable structure conditions placed upon the constraint function h , one may conclude analogous results to those presented here, albeit with the distance from a sphere in a possibly higher L^p norm [24, 25].

ACKNOWLEDGEMENTS

This work forms part of the author's PhD thesis under the support of an Australian Postgraduate Award. Part of this research was completed during two visits to the Freie Universität in Berlin. The first visit was under the support of the Deutscher Akademischer Austausch Dienst, and the second supported by the Research Group in Geometric Analysis at the Freie Universität under the supervision of Prof. Dr. Klaus Ecker. Part of this work was also completed during a visit to the Max Planck Institut in Gollm, under the supervision of Prof. Dr. Gerhard Huisken. He is grateful for their support and hospitality.

The author is also very grateful to the referee who carefully read the manuscript and suggested numerous improvements and simplifications to the arguments and exposition presented in this paper.

3. LOCAL ESTIMATES

This section is in two parts. In the first part we compute integral estimates which hold for any immersion, and use these with multiplicative Sobolev inequalities to show curvature estimates and a gap lemma. In the second part our primary focus is on proving estimates which will allow us to overcome our lack of a convenient gradient flow. There we utilise the evolution of certain integral quantities along the flow to conclude interior estimates and a lifespan theorem, similar to Theorems 3.5 and 1.2 in [14] and [15] respectively.

Throughout our analysis we will make extensive use of a cutoff function $\gamma = \tilde{\gamma} \circ f : M \rightarrow [0, 1]$ with compact support in M and $\tilde{\gamma} \in C_c^2(\mathbb{R}^3)$ which satisfies

$$(6) \quad \|\nabla \gamma\|_\infty \leq c_{\gamma 1}, \quad \|\nabla_{(2)} \gamma\|_\infty \leq c_{\gamma 2}(1 + |A|),$$

for some absolute constants $c_{\gamma 1}, c_{\gamma 2} < \infty$. Note γ automatically has compact support in M so long as f is a proper immersion. We frequently use the notation $\|\cdot\|_p$ for the L^p norm, and $\|\cdot\|_{p, [\gamma > 0]}$ for the L^p norm restricted to the set where a function γ is positive. We also use $P_i^j(T)$ to denote a sum of contractions of i copies of T with j covariant derivatives in each contraction.

Since we feel this argument is quite robust, we will use the convention $F = \partial_t^\perp f = \Delta H$ to indicate how one may proceed for other flows. We begin the first part of this section by noting some easy consequences of results from [14].

Proposition 4. *For an immersed hypersurface $f : M^n \rightarrow \mathbb{R}^{n+1}$, γ as in (6) and $s \geq 4$ we have*

$$\int_M |\nabla_{(2)} A|^2 \gamma^s d\mu + \int_M |A|^2 |\nabla A|^2 \gamma^s d\mu + \int_M |A|^4 |A^o|^2 \gamma^s d\mu$$

$$\leq c \int_M |F|^2 \gamma^s d\mu + c(c_{\gamma_1})^4 \int_{[\gamma>0]} |A|^2 d\mu + c_1 \int_M (|A^\circ|^2 |\nabla A^\circ|^2 + |A^\circ|^6) \gamma^s d\mu,$$

where c and c_1 are functions of n and s only.

Proof. Combine the obvious estimate

$$(7) \quad \int_M H^2 |A^\circ|^4 \gamma^s d\mu \leq \delta \int_M H^4 |A^\circ|^2 \gamma^s d\mu + \frac{1}{4\delta} \int_M |A^\circ|^6 \gamma^s d\mu,$$

valid for any $\delta > 0$, and Proposition 2.4 from [14]. \square \square

The next four lemmas are generalisations of results from [14, 15] to at least include the case $n = 3$. We include the first three with a view toward future applications, while the fourth is instrumental in proving a gap lemma for $n = 3$. Each proof relies in a crucial way on the exponent in the Michael-Simon Sobolev inequality [19], and for $n \geq 4$ we have not been able to obtain a statement as useful as (10). We will omit the proof of the first three lemmas as we feel they are quite straightforward.

Lemma 5. *Let $f : M^n \rightarrow \mathbb{R}^{n+1}$ be an immersed hypersurface. For $u \in C_c^1(M)$, $n < p \leq \infty$, $0 \leq \beta \leq \infty$ and $0 < \alpha \leq 1$ where $\frac{1}{\alpha} = (\frac{1}{n} - \frac{1}{p})\beta + 1$ we have*

$$\|u\|_\infty \leq c \|u\|_\beta^{1-\alpha} (\|\nabla u\|_p + \|Hu\|_p)^\alpha,$$

where $c = c(n, p, \beta)$.

Lemma 6. *Suppose $n \in \{2, 3\}$ and let γ be as in (6). Then for any tensor T on M^n and $s \geq 2$ there is a $c = c(n, s)$ such that*

$$\|T\gamma^s\|_\infty^4 \leq c \|T\gamma^s\|_2^{4-n} \left[\left(\int_M |\nabla_{(2)} T|^2 \gamma^{2s} d\mu \right)^{\frac{n}{2}} + \left(\int_M H^4 |T|^2 \gamma^{2s} d\mu \right)^{\frac{n}{2}} + (c_{\gamma_1})^{2n} \left(\int_{[\gamma>0]} |T|^2 d\mu \right)^{\frac{n}{2}} \right].$$

Lemma 7. *Let $n \in \{2, 3\}$. Then for any tensor T on $f : M^n \rightarrow \mathbb{R}^{n+1}$ and γ as in (6) there is a $c = c(c_{\gamma_1}, n)$ such that*

$$\|T\|_{\infty, [\gamma=1]}^4 \leq c \|T\|_{2, [\gamma>0]}^{4-n} (\|\nabla_{(2)} T\|_{2, [\gamma>0]}^n + \|TA^2\|_{2, [\gamma>0]}^n + \|T\|_{2, [\gamma>0]}^n).$$

Lemma 8. *Suppose γ is as in (6). Then for an immersion $f : M^2 \rightarrow \mathbb{R}^3$ and $s \geq 4$,*

$$(8) \quad \begin{aligned} & \int_M |A^\circ|^6 \gamma^s d\mu + \int_M |A^\circ|^2 |\nabla A^\circ|^2 \gamma^s d\mu \\ & \leq c_2 \int_{[\gamma>0]} |A^\circ|^2 d\mu \int_M (|\nabla_{(2)} A^\circ|^2 + |A|^2 |\nabla A^\circ|^2 + |A|^2 |A^\circ|^4) \gamma^s d\mu \\ & \quad + c(c_{\gamma_1})^4 \left(\int_{[\gamma>0]} |A^\circ|^2 d\mu \right)^2, \text{ and} \end{aligned}$$

$$(9) \quad \begin{aligned} & \int_M |A|^6 \gamma^s d\mu + \int_M |A|^2 |\nabla A|^2 \gamma^s d\mu \\ & \leq c \int_{[\gamma>0]} |A|^2 d\mu \int_M (|\nabla_{(2)} A|^2 + |A|^6) \gamma^s d\mu + c(c_{\gamma_1})^4 \left(\int_{[\gamma>0]} |A|^2 d\mu \right)^2, \end{aligned}$$

and for an immersion $f : M^3 \rightarrow \mathbb{R}^4$ and $s \geq 6$,

$$\begin{aligned}
& \int_M |A^\circ|^6 \gamma^s d\mu + \int_M |A^\circ|^2 |\nabla A^\circ|^2 \gamma^s d\mu \\
& \leq \delta \int_M |\nabla_{(2)} A^\circ|^2 \gamma^s d\mu + c_3 \|A^\circ\|_{3, [\gamma > 0]}^{\frac{3}{2}} \int_M (|\nabla_{(2)} A^\circ|^2 + |A|^2 |A^\circ|^4 + |A^\circ|^6) \gamma^s d\mu \\
(10) \quad & + c(c_{\gamma 1})^3 (\|A^\circ\|_{3, [\gamma > 0]}^3 + \|A^\circ\|_{3, [\gamma > 0]}^{\frac{9}{2}}),
\end{aligned}$$

where $\delta > 0$ and $c = c(s, n)$.

Proof. As the $n = 2$ case is Lemma 2.5 in [14] and Lemma 4.2 in [15], we only prove the $n = 3$ case. First, using integration by parts we estimate

$$(11) \quad \int |\nabla A^\circ|^3 \gamma^s d\mu \leq \delta \int_M |\nabla_{(2)} A^\circ|^2 \gamma^s d\mu + c_\delta \int_M |A^\circ|^6 \gamma^s d\mu + c(c_{\gamma 1})^3 \int_{[\gamma > 0]} |A^\circ|^3 d\mu,$$

for any $\delta > 0$. Now we use the Michael-Simon Sobolev inequality [19] with $u = |A^\circ|^4 \gamma^{2s/3}$ to estimate

$$\begin{aligned}
& \left(\int_M |A^\circ|^6 \gamma^s d\mu \right)^{\frac{2}{3}} \leq c \int_M |\nabla (|A^\circ|^4 \gamma^{2s/3})| d\mu + c \int_M |H| \cdot |A^\circ|^4 \gamma^{2s/3} d\mu \\
& \leq c \int_M |\nabla A^\circ|^2 |A^\circ| \gamma^s d\mu + \left(\int_M |A^\circ|^6 \gamma^s d\mu \right)^{\frac{2}{3}} \left(\int_{[\gamma > 0]} |A^\circ|^3 \right)^{\frac{1}{3}} \\
& \quad + c \left(\int_{[\gamma > 0]} |A^\circ|^3 d\mu \right)^{\frac{1}{3}} \left(\int_M |A|^2 |A^\circ|^4 \gamma^s d\mu \right)^{\frac{2}{3}} + c(c_{\gamma 1})^2 \|A^\circ\|_{3, [\gamma > 0]}^3.
\end{aligned}$$

Note that we needed $s \geq 6$. Estimating the above and then combining the result with (11) gives

$$\begin{aligned}
\int_M |A^\circ|^6 \gamma^s d\mu & \leq c \|A^\circ\|_{3, [\gamma > 0]}^{\frac{3}{2}} \int_M |\nabla A^\circ|^3 \gamma^s d\mu + c(c_{\gamma 1})^3 \|A^\circ\|_{3, [\gamma > 0]}^{\frac{9}{2}} \\
& \quad + c \|A^\circ\|_{3, [\gamma > 0]}^{\frac{3}{2}} \int_M (|A^\circ|^6 + |A|^2 |A^\circ|^4) \gamma^s d\mu \\
& \leq c \|A^\circ\|_{3, [\gamma > 0]}^{\frac{3}{2}} \int_M (|\nabla_{(2)} A^\circ|^2 + |A|^2 |A^\circ|^4 + |A^\circ|^6) \gamma^s d\mu \\
& \quad + c(c_{\gamma 1})^3 \|A^\circ\|_{3, [\gamma > 0]}^{\frac{9}{2}}.
\end{aligned}$$

This estimates the first term. For the second, simply estimate

$$\int_M |A^\circ|^2 |\nabla A^\circ|^2 \gamma^s d\mu \leq c \int_M |A^\circ|^6 \gamma^s d\mu + c \int_M |\nabla A^\circ|^3 \gamma^s d\mu$$

and again use (11). This estimates the second term, and combining the two estimates above finishes the proof. \square \square

We can now prove our main estimate for this part.

Proposition 9. *Suppose $n \in \{2, 3\}$, γ is as in (6) and $f : M^n \rightarrow \mathbb{R}^{n+1}$ is an immersion. Then there exists an $\epsilon_0 > 0$ such that if $\|A^\circ\|_{n, [\gamma > 0]}^n < \epsilon_0$ we have*

$$\begin{aligned} & \int_M (|\nabla_{(2)}A|^2 + |A|^2|\nabla A|^2 + |A|^4|A^\circ|^2)\gamma^6 d\mu \\ & \leq c \int_M |F|^2\gamma^6 d\mu + c(c_{\gamma 1})^4 \int_{[\gamma > 0]} |A|^2 d\mu + c(c_{\gamma 1})^{6-n} \|A^\circ\|_{n, [\gamma > 0]}^n, \end{aligned}$$

where c depends on n only.

Proof. First note that

$$(12) \quad \|A^\circ\|_{2, [\gamma > 0]}^4 \leq \epsilon_0 \|A^\circ\|_{2, [\gamma > 0]}^2, \quad \text{and} \quad \|A^\circ\|_{3, [\gamma > 0]}^{\frac{9}{2}} \leq \epsilon_0^{\frac{1}{2}} \|A^\circ\|_{3, [\gamma > 0]}^3,$$

for $n = 2$ and $n = 3$ respectively. Combining Proposition 4 with (8) for $n = 2$, (10) for $n = 3$, gives

$$\begin{aligned} & \int_M (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2)\gamma^s d\mu \\ & \leq c \int_M |F|^2\gamma^s d\mu + c(c_{\gamma 1})^4 \|A\|_{2, [\gamma > 0]}^2 + c_1\delta(n-2) \int_M |\nabla_{(2)}A|^2\gamma^s d\mu \\ & \quad + c_1c_n \|A^\circ\|_{n, [\gamma > 0]}^{\frac{n}{n-1}} \int_M (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2)\gamma^s d\mu \\ & \quad + c(n-2)(c_{\gamma 1})^3 (\|A^\circ\|_{3, [\gamma > 0]}^3 + \|A^\circ\|_{3, [\gamma > 0]}^{\frac{9}{2}}) \\ & \quad + c(3-n)(c_{\gamma 1})^4 \|A^\circ\|_{2, [\gamma > 0]}^4, \end{aligned}$$

for any $\delta > 0$. Absorbing and using (12),

$$\begin{aligned} & (1 - c_1c_n\epsilon_0^{\frac{1}{n-1}} - c_1\delta(n-2)) \int_M (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2)\gamma^6 d\mu \\ & \leq c \int_M |F|^2\gamma^6 d\mu + c(c_{\gamma 1})^4 \|A\|_{2, [\gamma > 0]}^2 + c(c_{\gamma 1})^{6-n} \|A^\circ\|_{n, [\gamma > 0]}^n. \end{aligned}$$

Therefore, for δ, ϵ_0 sufficiently small we may conclude the result. \square \square

We can now give our main theorem for this part.

Theorem 10 (Gap Lemma). *Suppose $n \in \{2, 3\}$ and $f : M^n \rightarrow \mathbb{R}^{n+1}$ is a properly immersed surface with $F = \Delta H \equiv 0$. Then if*

$$\int_M |A^\circ|^n d\mu < \epsilon_0, \quad \text{and} \quad \liminf_{\rho \rightarrow \infty} \frac{1}{\rho^4} \int_{f^{-1}(B_\rho(0))} |A|^2 d\mu = 0,$$

f is an embedded plane or sphere.

Proof. We set the cutoff function γ to be such that $\gamma(p) = \varphi\left(\frac{1}{\rho}|f(p)|\right)$, where $\varphi \in C^1(\mathbb{R})$ and $\chi_{B_{1/2}(0)} \leq \varphi \leq \chi_{B_1(0)}$. With this choice $c_{\gamma 1} = \frac{c}{\rho}$. Recall that in our estimates we do not use the second derivative of γ . Taking $\rho \rightarrow \infty$ in Proposition 9 gives

$$\int_M |\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2 d\mu \leq c \liminf_{\rho \rightarrow \infty} \frac{1}{\rho^4} \|A\|_{2, f^{-1}(B_\rho(0))}^2,$$

where the terms involving A° vanished since $\|A^\circ\|_2^n < \epsilon_0 < \infty$. Therefore $|A| = 0$, $|A^\circ| = 0$, or both, which implies that f maps into a round sphere or plane $S \subset \mathbb{R}^{n+1}$, and since f is complete the map $f : M \rightarrow S$ is a global isometry. \square \square

We now begin the second part of this section, changing focus by proving estimates which rely on the parabolic nature of the evolution equation (1). To begin, we state the following elementary evolution equations, whose proof is standard [12].

Lemma 11. *For a surface diffusion flow $f : M^n \times [0, T] \rightarrow \mathbb{R}^{n+1}$ the following equations hold:*

$$\begin{aligned} \frac{\partial}{\partial t} g_{ij} &= 2FA_{ij}, & \frac{\partial}{\partial t} g^{ij} &= -2FA^{ij}, & \frac{\partial}{\partial t} d\mu &= (HF)d\mu, \\ \frac{\partial}{\partial t} \nu &= -\nabla F, & \frac{\partial}{\partial t} A_{ij} &= -\nabla_{ij} F + FA_i^p A_{pj}, \\ \frac{\partial}{\partial t} A_{ij}^o &= -S^o(\nabla_{(2)} F) + F[A_i^p A_{pj} + \frac{1}{n} g_{ij} |A|^2 - \frac{2}{n} H A_{ij}], & \text{and} \\ \frac{\partial}{\partial t} \nabla_{(k)} A &= -\Delta^2 \nabla_{(k)} A + P_3^{k+2}(A), & k \in \mathbb{N}_0, \end{aligned}$$

where $S^o(T)$ is the symmetric tracefree part of a bilinear form T .

Using the above evolution equations, and the previous estimates, we may now prove that the Willmore energy is nonincreasing for surface diffusion flows satisfying (4).

Proposition 12. *Suppose $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ is a surface diffusion flow satisfying (4). Then*

$$\frac{d}{dt} \int_M |A^o|^2 d\mu = \frac{1}{2} \frac{d}{dt} \int_M |H|^2 d\mu \leq -\frac{1}{4} \int_M |\Delta H|^2 d\mu,$$

and $f(\cdot, t)$ is a family of embeddings.

Proof. Proposition 9 with $\gamma \equiv 1$ gives

$$\int_M |\nabla_{(2)} A|^2 + |A|^2 |\nabla A|^2 + |A|^4 |A^o|^2 d\mu \leq c \int_M |\Delta H|^2 d\mu.$$

Combining this with (7) and (8) we obtain

$$\begin{aligned} \int_M |A^o|^4 H^2 d\mu &\leq c\delta \int_M |\Delta H|^2 d\mu + c_\delta \int_M |A^o|^6 d\mu \\ (13) \qquad \qquad \qquad &\leq c(\delta + c_\delta \epsilon_0) \int_M |\Delta H|^2 d\mu. \end{aligned}$$

Using Gauss-Bonnet, Lemma 11 and (13) we compute

$$\begin{aligned} \frac{d}{dt} \int_M |A^o|^2 d\mu &= \frac{1}{2} \frac{d}{dt} \int_M |H|^2 d\mu \\ &= - \int_M (\Delta H + H |A^o|^2) (\Delta H) d\mu \\ &\leq -\frac{1}{2} \int_M |\Delta H|^2 d\mu + c \int_M |A^o|^4 H^2 d\mu \\ &\leq -\frac{1}{4} \int_M |\Delta H|^2 d\mu, \end{aligned}$$

for δ, ϵ_0 sufficiently small. We have thus shown, for all t ,

$$\frac{1}{4} \int_M H^2 d\mu < 8\pi,$$

and it follows from Theorem 6 in [16] that each $f(\cdot, t)$ is an embedding. $\square \quad \square$

Lemma 13. *Let γ be as in (6). The following equalities hold for a surface diffusion flow $f : M^n \times [0, T) \rightarrow \mathbb{R}^{n+1}$:*

$$\begin{aligned}
& \frac{d}{dt} \int_M \frac{1}{2} H^2 \gamma^s d\mu + \int_M |F|^2 \gamma^s d\mu = \frac{1}{2} \int_M H^2 (\partial_t \gamma^s) d\mu - \int_M FH |A^\circ|^2 \gamma^s d\mu \\
& \quad + \int_M H \langle \nabla F, \nabla \gamma^s \rangle - F \langle \nabla H, \nabla \gamma^s \rangle d\mu, \text{ and} \\
& \frac{d}{dt} \int_M |A^\circ|^2 \gamma^s d\mu + \int_M |F|^2 \gamma^s d\mu = \int_M |A^\circ|^2 (\partial_t \gamma^s) d\mu + \int_M FH |A^\circ|^2 \gamma^s d\mu \\
& \quad + 2 \int_M H \langle A^\circ, \nabla F \nabla \gamma^s \rangle + F \langle \nabla^* A^\circ, \nabla \gamma^s \rangle d\mu \\
(14) \quad & \quad - 2 \int_M F (A^\circ)_j^i (A^\circ)_k^j (A^\circ)_i^k \gamma^s d\mu.
\end{aligned}$$

Proof. Using Lemma 11,

$$\begin{aligned}
\frac{d}{dt} \int_M \frac{1}{2} H^2 \gamma^s d\mu &= \int_M H (-\Delta F - |A|^2 F) \gamma^s d\mu + \int_M \frac{1}{2} H^2 (\partial_t \gamma^s) d\mu + \int_M \frac{1}{2} H^3 F \gamma^s d\mu \\
&= - \int_M H (\Delta F) \gamma^s d\mu - \int_M FH |A^\circ|^2 \gamma^s d\mu + \frac{1}{2} \int_M H^2 (\partial_t \gamma^s) d\mu.
\end{aligned}$$

Integrating by parts twice,

$$\begin{aligned}
\frac{d}{dt} \int_M \frac{1}{2} H^2 \gamma^s d\mu + \int_M F^2 \gamma^s d\mu &= \frac{1}{2} \int_M H^2 (\partial_t \gamma^s) d\mu - \int_M FH |A^\circ|^2 \gamma^s d\mu \\
& \quad + \int_M H \langle \nabla F, \nabla \gamma^s \rangle - F \langle \nabla H, \nabla \gamma^s \rangle d\mu.
\end{aligned}$$

This proves the first statement. For the second, first compute

$$\begin{aligned}
\frac{d}{dt} \int_M |A^\circ|^2 \gamma^s d\mu &= 2 \int_M \left\langle A_{ij}^\circ, -S_{ij}^\circ (\nabla_{(2)} F) + F [A_i^p A_{pj} + \frac{1}{n} g_{ij} |A|^2 - \frac{2}{n} H A_{ij}] \right\rangle \gamma^s d\mu \\
& \quad + \int_M |A^\circ|^2 H F \gamma^s d\mu + \int_M |A^\circ|^2 (\partial_t \gamma^s) d\mu \\
&= -2 \int_M \langle A^\circ, \nabla_{(2)} F \rangle \gamma^s d\mu - 2 \int_M F (A^\circ)_j^i (A^\circ)_k^j (A^\circ)_i^k \gamma^s d\mu \\
& \quad + \int_M |A^\circ|^2 H F \gamma^s d\mu + \int_M |A^\circ|^2 (\partial_t \gamma^s) d\mu,
\end{aligned}$$

since $\langle A_{ij}^\circ, A_i^p A_{pj} - \frac{2}{n} H A_{ij} \rangle = (A^\circ)_{ij} (A^\circ)_k^i (A^\circ)^{kj}$. Using now $\nabla_{(2)}^* A^\circ = \frac{n-1}{n} \Delta H$ and integration by parts twice gives the second statement, and so we are finished. \square

Considering for the moment the tracefree curvature estimate only, we tailor the right hand side of the previous identity to be compatible with the multiplicative Sobolev inequalities of Lemma 8.

Lemma 14. *Let γ be as in (6), $s \geq 4$, and $\delta_1, \delta_2, \delta_3$ be fixed positive numbers. The following inequality holds for a surface diffusion flow $f : M^n \times [0, T) \rightarrow \mathbb{R}^{n+1}$:*

$$\begin{aligned}
\frac{d}{dt} \int_M |A^\circ|^2 \gamma^s d\mu + (1 - \delta_1) \int_M F^2 \gamma^s d\mu &\leq \delta_2 \int_M |\nabla A^\circ|^2 H^2 \gamma^s d\mu \\
& \quad + \delta_3 \int_M |A|^4 |A^\circ|^2 \gamma^s d\mu + c \int_M (|A^\circ|^6 + |\nabla A^\circ|^2 |A^\circ|^2) \gamma^s d\mu
\end{aligned}$$

$$+ c[(c_{\gamma 2})^2 + (c_{\gamma 1})^4] \int_M |A^\circ|^2 \gamma^{s-4} d\mu.$$

where $c = c(s, n, \delta_i)$.

Proof. We first compute, using integration by parts and the definition of γ ,

$$\begin{aligned} & 2 \int_M (A^\circ)_{ij} (\nabla^i \gamma^s) (\nabla^j F) d\mu + 2 \int_M F \langle \nabla^* A^\circ, \nabla \gamma^s \rangle d\mu = -2 \int_M F \langle A^\circ, \nabla_{(2)} \gamma^s \rangle d\mu \\ & = -2s \int_M F \left(\gamma \langle A^\circ, (D^2 \tilde{\gamma} \circ f) g \rangle + (D \tilde{\gamma} \circ f) A \right) + (s-1) \langle A^\circ, \nabla \gamma \nabla \gamma \rangle \gamma^{s-2} d\mu. \end{aligned}$$

Note that $\langle A^\circ, (D^2 \tilde{\gamma} \circ f) g \rangle = \sum_{i,j=1}^n A^\circ_{ij} (D_{ij} \tilde{\gamma}) g_{ij}$ is not in general zero, as each term in the sum is scaled by the second derivatives of $\tilde{\gamma}$. Continuing,

$$\begin{aligned} & 2 \int_M (A^\circ)_{ij} (\nabla^i \gamma^s) (\nabla^j F) d\mu + 2 \int_M F \langle \nabla^* A^\circ, \nabla \gamma^s \rangle d\mu \\ & \leq c \int_M |F| (|A^\circ| (c_{\gamma 2}) + |A^\circ|^2 (c_{\gamma 1})) \gamma^{s-1} d\mu + c \int_M |F| \cdot |A^\circ| (c_{\gamma 1})^2 \gamma^{s-2} d\mu \\ & \leq \frac{\delta_1}{2} \int_M |F|^2 \gamma^s d\mu + \int_M |A^\circ|^6 \gamma^s d\mu + c \delta_1^{-2} [(c_{\gamma 2})^2 + (c_{\gamma 1})^4] \int_M |A^\circ|^2 \gamma^{s-4} d\mu, \end{aligned}$$

We now estimate the terms $\int_M FH|A^\circ|^2 \gamma^s d\mu$ and $\int_M F(A^\circ)^i_j (A^\circ)^j_k (A^\circ)^k_i \gamma^s d\mu$ from Lemma 13. Using integration by parts and $\nabla_{(2)}^* A^\circ = \frac{n-1}{n} \Delta H$,

$$\begin{aligned} \int_M FH|A^\circ|^2 \gamma^s d\mu & = - \int_M |\nabla H|^2 |A^\circ|^2 \gamma^s d\mu - s \int_M H |A^\circ|^2 \langle \nabla H, \nabla \gamma \rangle \gamma^{s-1} d\mu \\ & \quad - 2 \int_M H |A^\circ| \langle \nabla H, \nabla |A^\circ| \rangle \gamma^s d\mu \\ & \leq c_{\delta_3} \int_M |\nabla A^\circ|^2 |A^\circ|^2 \gamma^s d\mu + \delta_3 \int_M |\nabla A^\circ|^2 H^2 \gamma^s d\mu \\ & \quad + c_{\delta_2} (c_{\gamma 1})^4 \int_M |A^\circ|^2 \gamma^{s-4} d\mu + \delta_2 \int_M |A|^4 |A^\circ|^2 \gamma^s d\mu. \end{aligned}$$

This estimates the first integral. The second is easily estimated by

$$\int_M F(A^\circ * A^\circ * A^\circ) \gamma^s d\mu \leq \frac{\delta_1}{4} \int_M F^2 \gamma^s d\mu + c_{\delta_1} \int_M |A^\circ|^6 \gamma^s d\mu.$$

Finally we estimate the time derivative of γ as follows.

$$\begin{aligned} \int_M |A^\circ|^2 (\partial_t \gamma^s) d\mu & = s \int_M |A^\circ|^2 (\gamma^{s-1} \langle D \tilde{\gamma}, \nu \rangle F) d\mu \\ & \leq \frac{\delta_1}{4} \int_M F^2 \gamma^s d\mu + c_{\delta_1} (c_{\tilde{\gamma} 1})^4 \int_M |A^\circ|^2 \gamma^{s-4} d\mu + \int_M |A^\circ|^6 d\mu. \end{aligned}$$

Combining these inequalities with Lemma 13 and absorbing finishes the proof. \square

The following proposition demonstrates some local control of the L^2 norm of the tracefree curvature.

Proposition 15. *Suppose $f : M^2 \times [0, T^*] \rightarrow \mathbb{R}^3$ is a surface diffusion flow where for all $x \in \mathbb{R}^3$ there exists $\sigma < \infty$ such that $\|A\|_{2, f^{-1}(B_\rho(x))}^2 \leq \sigma$. Then there exist constants $\epsilon_1 > 0$ and $c_4 = c_4(\epsilon_0) > 0$ such that if*

$$(15) \quad \sup_{x \in \mathbb{R}^3} \int_{f^{-1}(B_\rho(x))} |A^\circ|^2 d\mu \Big|_{t=0} < \epsilon_1$$

then at any time $0 \leq t < t_1 = \min\{c_4\sigma^{-1}\rho^4, T^*\}$ and for any $x \in \mathbb{R}^3$ we have

$$(16) \quad \int_{f^{-1}(B_\rho(x))} |A^\circ|^2 d\mu + \frac{1}{2} \int_0^t \int_{f^{-1}(B_\rho(x))} (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2) d\mu d\tau \leq c[\epsilon_1 + \sigma\rho^{-4}t].$$

Proof. We first establish an estimate for the tracefree curvature. Assuming that (15) is satisfied on $[0, t]$, $t < t_1$, we combine Proposition 9 with the multiplicative Sobolev inequality (8) and Lemma 14 to obtain

$$\begin{aligned} \frac{d}{dt} \int_M |A^\circ|^2 \gamma^s d\mu + c(1 - \sum \delta_i - c_2\epsilon_0) \int_M (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2) \gamma^s d\mu \\ \leq \frac{c}{\rho^4} \|A^\circ\|_{2, [\gamma>0]}^4 + \frac{c}{\rho^4} \|A\|_{2, [\gamma>0]}^2 \leq \frac{c}{\rho^4} \|A\|_{2, [\gamma>0]}^2. \end{aligned}$$

where $\epsilon_0 > 0$ is as in Proposition 9. Choosing δ_i small enough and requiring ϵ_0 to be such that $c(1 - \sum_i \delta_i - c_2\epsilon_0) = \frac{1}{2}$ we have shown

$$(17) \quad \frac{d}{dt} \int_M |A^\circ|^2 \gamma^s d\mu + \frac{1}{2} \int_M (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2) \gamma^s d\mu \leq \frac{c}{\rho^4} \|A\|_{2, [\gamma>0]}^2.$$

Now integrate (17) to obtain

$$(18) \quad \int_{f^{-1}(B_{\rho/2}(x))} |A^\circ|^2 d\mu + \frac{1}{2} \int_0^t \int_{f^{-1}(B_{\rho/2}(x))} (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2) d\mu d\tau \leq \epsilon_1 + c\sigma\rho^{-4}t.$$

Motivated by the fact that 2^4 balls $B_{\rho/2}$ can be used to cover a ball B_ρ , we require $\epsilon_1 \leq \frac{\epsilon_0}{4 \cdot 2^4}$. Since $0 < t \leq \frac{\epsilon_0}{c_4 \cdot 2^4} \rho^4 \sigma^{-1} = c_4 \sigma^{-1} \rho^4$, we use (18) and a covering argument to derive

$$\begin{aligned} \int_{f^{-1}(B_\rho(x))} |A^\circ|^2 d\mu + \frac{1}{2} \int_0^t \int_{f^{-1}(B_\rho(x))} (|\nabla_{(2)}A|^2 + |\nabla A|^2|A|^2 + |A|^4|A^\circ|^2) d\mu d\tau \\ \leq 2^4 \left(\epsilon_0 \frac{1}{4 \cdot 2^4} + \epsilon_0 \frac{1}{4 \cdot 2^4} \right) \leq \frac{\epsilon_0}{2}. \end{aligned}$$

Therefore the smallness assumption of Proposition 9 holds up to time $t = t_1$ and (16) follows. \square

For the interested reader we also state an easy corollary.

Corollary 16. *Under the assumptions of Proposition 15 above,*

$$\begin{aligned} \int_0^t \|A^\circ\|_{\infty, f^{-1}(B_\rho(x_1))}^4 d\tau \leq c[\epsilon_1 + \sigma\rho^{-4}t], \text{ and} \\ \int_{f^{-1}(B_{\rho'/2}(x_1))} |A|^2 d\mu \Big|_{t=\tau} \leq \int_{f^{-1}(B_{\rho'}(x_1))} |A|^2 d\mu \Big|_{t=0} + c(\sigma(\rho')^{-4})\tau, \end{aligned}$$

where $0 < \rho' \leq \rho$ and $\tau \leq \min\{c_4\sigma^{-1}(\rho')^4, T^*\}$.

Proof. To obtain the first estimate, combine (16) with Lemma 7 and use a covering argument. For the second, first note that an argument completely analogous to that used to obtain (17) also gives

$$\frac{d}{dt} \int_M H^2 \gamma^s d\mu + \frac{1}{2} \int_M (|\nabla_{(2)} A|^2 + |\nabla A|^2 |A|^2 + |A|^4 |A^o|^2) \gamma^s d\mu \leq \frac{c}{\rho^4} \|A\|_{2, [\gamma > 0]}^2.$$

Now the second statement follows from integrating this and combining with the estimate of Proposition 15. \square \square

We now turn to controlling the higher derivatives of curvature. Using Lemma 11, one may derive the following local integral estimate. The proof is long but somewhat standard; see [15] for the case of Willmore flow. One computes the time derivative of $\|(\nabla_{(k)} A) \gamma^s\|_2^2$, integrates by parts, estimates the result, interpolates, and absorbs higher derivative terms on the left.

Proposition 17. *Let $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ be a surface diffusion flow and γ as in (6). Then for a fixed $\theta > 0$ and $s \geq 2k + 4$,*

$$\begin{aligned} & \frac{d}{dt} \int_M |\nabla_{(k)} A|^2 \gamma^s d\mu + (2 - \theta) \int_M |\nabla_{(k+2)} A|^2 \gamma^s d\mu \\ & \leq c \int_M |A|^2 \gamma^{s-4-2k} d\mu + c \int_M ([P_3^{k+2}(A) + P_5^k(A)] * \nabla_{(k)} A) \gamma^s d\mu, \end{aligned}$$

where $c = c(c_{\gamma_1}, c_{\gamma_2}, s, k, \theta)$.

Note that in the above local smallness of curvature is not assumed. Using this we obtain the following.

Proposition 18. *Suppose $f : M^2 \times [0, T^*] \rightarrow \mathbb{R}^3$ is a surface diffusion flow and γ is as in (6). Then there is an $\epsilon_0 > 0$ such that if*

$$\epsilon = \sup_{[0, T^*]} \int_{[\gamma > 0]} |A|^2 d\mu \leq \epsilon_0$$

then for any $t \in [0, T^*]$ we have

$$(19) \quad \begin{aligned} & \int_{[\gamma=1]} |A|^2 d\mu + \int_0^t \int_{[\gamma=1]} (|\nabla_{(2)} A|^2 + |A|^2 |\nabla A|^2 + |A|^6) d\mu d\tau \\ & \leq \int_{[\gamma > 0]} |A|^2 d\mu \Big|_{t=0} + c\epsilon t, \end{aligned}$$

where $c = c(c_{\gamma_1}, c_{\gamma_2})$.

The idea of the proof is to integrate Proposition 17, and then use the multiplicative Sobolev inequality (9). This will introduce a multiplicative factor of $\|A\|_{2, [\gamma > 0]}^2$ in front of several integrals, which we can then absorb on the left. Since this is similar to [15], we omit the proof.

We will need the below proposition to obtain interior estimates. This is proved using Proposition 17, interpolating the P -terms using Corollary 5.5 from [15] and then absorbing.

Proposition 19. *Suppose $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ is a surface diffusion flow and γ a localisation function as in (6). Then, for $s \geq 2k + 4$ the following estimate holds:*

$$(20) \quad \begin{aligned} & \frac{d}{dt} \int_M |\nabla_{(k)} A|^2 \gamma^s d\mu + \int_M |\nabla_{(k+2)} A|^2 \gamma^s d\mu \\ & \leq c \|A\|_{\infty, [\gamma > 0]}^4 \int_M \|\nabla_{(k)} A\|^2 \gamma^s d\mu + c \|A\|_{2, [\gamma > 0]}^2 (1 + \|A\|_{\infty, [\gamma > 0]}^4). \end{aligned}$$

We next state local interior estimates for (1), which hold where and when curvature is locally small. Once the estimates (19) and (20) are established, the proof follows similarly to Theorem 3.5 in [14]. Briefly, one utilises cutoff functions in time as well as the familiar spatial localisation functions γ to obtain integral estimates of Gronwall type, which in turn gives the estimates in L^2 . Then, using interpolation (such as in Lemma 7), one obtains the L^∞ estimates.

Theorem 20 (Interior estimates). *Suppose $f : M^2 \times (0, T^*] \rightarrow \mathbb{R}^3$ is a surface diffusion flow satisfying*

$$\sup_{0 < t \leq T^*} \int_{f^{-1}(B_\rho(x))} |A|^2 d\mu \leq \epsilon_0, \text{ for } T^* \leq c\rho^4.$$

Then for any $k \in \mathbb{N}_0$ we have at time $t \in (0, T^]$ the estimates*

$$\begin{aligned} \|\nabla_{(k)} A\|_{2, f^{-1}(B_{\rho/2}(x))} & \leq c_k \sqrt{\epsilon_0} t^{-\frac{k}{4}} \\ \|\nabla_{(k)} A\|_{\infty, f^{-1}(B_{\rho/2}(x))} & \leq c_k \sqrt{\epsilon_0} t^{-\frac{k+1}{4}}, \end{aligned}$$

where $c_k = c_k(k, \rho, T^, \|\nabla_{(k)} A\|_{2, f^{-1}(B_\rho(x_0))}|_{t=0})$.*

Armed with these estimates, we may conclude the following concentration-compactness alternative. The proof is similar to that of Theorem 1.2 in [15]: one defines a time parameter t_0 and corresponding interval $[0, t_0]$, which is the time during which the curvature is locally small in L^2 . Using a continuity argument, one shows that either $t_0 = c_0\rho^4$, in which case one can conclude the theorem, or that $t_0 = T$. For the latter case, we then use the interior estimates and a tiling argument such as can be found in [10] or [11] to show that $T = \infty$, and we again conclude the theorem.

Theorem 21 (Lifespan Theorem). *Suppose $f : M^2 \times [0, T) \rightarrow \mathbb{R}^3$ is a surface diffusion flow with f_0 smooth. Then there are constants $\rho, \epsilon_0, c_0 > 0$ such that if ρ is chosen with*

$$\int_{f^{-1}(B_\rho(x))} |A|^2 d\mu \Big|_{t=0} \leq \epsilon_0 \quad \text{for any } x \in \mathbb{R}^3,$$

then the maximal time T satisfies

$$T \geq c_0\rho^4,$$

and we have the estimate

$$\int_{f^{-1}(B_\rho(x))} |A|^2 d\mu \leq c\epsilon_0 \quad \text{for } 0 \leq t \leq c_0\rho^4.$$

22. BLOWUP ANALYSIS

This section is split into several parts. First we state the required compactness theorem, and detail the construction of the blowup. Our primary concern is to demonstrate that under suitable conditions the blowup is a stationary solution of surface diffusion flow. Although we do not enjoy an explicit gradient flow of curvature, Proposition 12 tells us that surface diffusion flow is ‘almost’ a gradient flow of curvature, in the sense that if the tracefree curvature is initially globally small enough, this condition is preserved and in fact the L^2 norm of the curvature is nonincreasing. Combining this with our estimates from Section 2, we are able to essentially follow [14]. For the convenience of the reader, we detail the entire argument.

Observe that by Theorem 21, $T < \infty$ implies that the curvature has concentrated in L^2 (see (22)). Assuming (4), we prove that the blowup at a time where the curvature has concentrated in L^2 is a stationary, noncompact, nonumbilic surface satisfying the assumptions of Theorem 10. We apply this in the next section, where we contradict finite maximal existence time with Theorem 10 to conclude long time existence. A further argument strengthens this result to exponential convergence to spheres, and this finishes the proof of Theorem 2.

Theorem 23 (Theorem 4.2, [14]). *Let $f_j : M_j \rightarrow \mathbb{R}^3$ be a sequence of proper immersions, where M_j is a surface without boundary. Let*

$$M_j(R) = \{p \in M_j : |f_j(p)| < R\}$$

and assume the bounds

$$\begin{aligned} \mu_j(M_j(R)) &\leq c(R) \quad \text{for any } R > 0, \\ \|\nabla_{(k)} A\|_{\infty, M_j} &\leq c(k) \quad \text{for any } k \in \mathbb{N}_0. \end{aligned}$$

Then there exists a proper immersion $\tilde{f} : \tilde{M} \rightarrow \mathbb{R}^3$, where \tilde{M} is again a surface without boundary, such that after passing to a subsequence we have a representation

$$f_j \circ \phi_j = \tilde{f} + u_j \quad \text{on } \tilde{M}(j) = \{p \in \tilde{M} : |\tilde{f}(p)| < j\}$$

with the following properties:

$$\begin{aligned} \phi_j : \tilde{M}(j) &\rightarrow U_j \subset M_j \text{ is diffeomorphic,} \\ M_j(R) &\subset U_j \text{ if } j \geq j(R), \\ u_j &\in C^\infty(\tilde{M}(j), \mathbb{R}^3) \text{ is normal along } \tilde{f}, \\ (21) \quad \|\tilde{\nabla}_{(k)} u_j\|_{\infty, \tilde{M}(j)} &\rightarrow 0 \text{ as } j \rightarrow \infty \text{ for any } k \in \mathbb{N}_0. \end{aligned}$$

The theorem says that on any ball $B_R(0)$ the immersion f_j can be written as a normal graph with small norm for j large over a limit immersion \tilde{f} , after suitably reparametrising with ϕ_j .

Let $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ be a surface diffusion flow defined on a closed surface M^2 , where $0 < T \leq \infty$. Define

$$\eta(r, t) = \sup_{x \in \mathbb{R}^3} \int_{f^{-1}(B_r(x))} |A|^2 d\mu.$$

Let r_j be an arbitrary decreasing sequence with $r_j \searrow 0$ and assume curvature concentrates in the sense that for each j ,

$$(22) \quad t_j = \inf\{t \geq 0 : \eta(r_j, t) > \epsilon_1\} < T,$$

where $\epsilon_1 = \epsilon_0 c_0$ and ϵ_0, c_0 are as in the Lifespan Theorem.

Lemma 24. *With the definitions above, we have*

$$\int_{f^{-1}(B_{r_j}(x))} |A|^2 d\mu \Big|_{t=t_j} \leq \epsilon_1 \text{ for any } x \in \mathbb{R}^3,$$

and

$$\int_{f^{-1}(\overline{B_{r_j}(x_j)})} |A|^2 d\mu \Big|_{t=t_j} \geq \epsilon_1 \text{ for some } x_j \in \mathbb{R}^3.$$

Proof. The first statement is a direct consequence of the definition of t_j . For the second, fix j and consider a sequence $\nu \rightarrow \infty$, $\nu > 1$. Consider times $t_{j+\nu-1} \searrow t_j$ and radii $r_{j+\nu-2} \nearrow r_j$. Then for each ν there exists an $x_{j+\nu-1}$ such that

$$\int_{f^{-1}(B_{r_{j+\nu-2}}(x_{j+\nu-1}))} |A|^2 d\mu \Big|_{t=t_{j+\nu-1}} \geq \epsilon_1.$$

Taking $\nu \rightarrow \infty$ in the above equation gives the second statement. \square \square

We now rescale f . Define immersions

$$f_j : M^2 \times [-r_j^{-4}t_j, r_j^{-4}(T-t_j)] \rightarrow \mathbb{R}^3, \quad f_j(p, t) = \frac{1}{r_j}(f(p, t_j + r_j^4 t) - x_j).$$

The sequence of immersions f_j can be thought of as ‘zooming in’ on the assumed curvature singularity at time T . Let $\eta_j(r, t)$ be η with respect to the immersion f_j . Then from Lemma 24 we have $\eta_j(1, t) \leq \epsilon_1$ for $t \leq 0$ and

$$\int_{f_j^{-1}(B_1(0))} |A|^2 d\mu \Big|_{t=0} \geq \epsilon_1.$$

The Lifespan Theorem implies $r_j^{-4}(T-t_j) \geq c_0$ for any j and also that

$$(23) \quad \eta_j(1, t) \leq \epsilon_0 \quad \text{for } 0 < t \leq c_0.$$

We apply the interior estimates on parabolic cylinders $B_1(x) \times (t-1, t]$ to obtain

$$(24) \quad \|\nabla_{(k)} A\|_{\infty, f_j} \leq c(k) \quad \text{for } -r_j^{-4}t_j + 1 \leq t \leq c_0.$$

Surface area is uniformly bounded, from above and below. Note that in light of Proposition 12 we could also follow [14] and use the local area estimate of Simon [22]. Using Theorem 23 with the sequence $f_j = f_j(\cdot, 0) : M^2 \rightarrow \mathbb{R}^3$ we obtain a limit immersion $\tilde{f}_0 : \tilde{M} \rightarrow \mathbb{R}^3$. Let $\phi_j : \tilde{M}(j) \rightarrow U_j \subset M^2$ be as in (21). Then the reparametrisation

$$(25) \quad f_j(\phi_j, \cdot) : \tilde{M}(j) \times [0, c_0] \rightarrow \mathbb{R}^3$$

is a surface diffusion flow with initial data

$$f_j(\phi_j, 0) = \tilde{f}_0 + u_j : \tilde{M}(j) \rightarrow \mathbb{R}^3.$$

The flows (25) satisfy the curvature bounds (24) and have initial data converging locally in C^k to the immersion $\tilde{f}_0 : M \rightarrow \mathbb{R}^3$. By converting the curvature bounds to partial derivative bounds in parabolic cylinders we obtain the locally smooth convergence

$$(26) \quad f_j(\phi_j, \cdot) \rightarrow \tilde{f},$$

where $\tilde{f} : \tilde{M} \times [0, c_0] \rightarrow \mathbb{R}^3$ is a surface diffusion flow with initial data \tilde{f}_0 .

Theorem 25. *Let $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ be a surface diffusion flow satisfying (4). Then the blowup \tilde{f} as constructed above is stationary.*

Proof. Noting the scale invariance of $\|A\|_2^2$, we use Proposition 12 to compute

$$\begin{aligned} \int_0^{c_0} \int_{\tilde{M}(j)} |\Delta H(f_j(\phi_j, t))|^2 d\mu_{f_j(\phi_j, \cdot)} dt &= \int_0^{c_0} \int_{U_j} |\Delta H_j|^2 d\mu_j dt \\ &\leq \int_M |A_j(0)|^2 d\mu_j - \int_M |A_j(c_0)|^2 d\mu_j \\ &= \int_M |A(t_j)|^2 d\mu - \int_M |A(t_j + r_j^4 c_0)|^2 d\mu, \end{aligned}$$

and this converges to zero as $j \rightarrow \infty$. Therefore $\Delta H(\tilde{f}) \equiv 0$ and the blowup \tilde{f} is stationary under surface diffusion flow. \square \square

Remark. *The above is strictly weaker than the corresponding result for Willmore flow from [14]. In particular, we need the tracefree curvature to be small at initial time to obtain a stationary blowup at a finite time curvature singularity, whereas for Willmore flow the corresponding blowup is always a stationary Willmore surface. Note that this means one may still obtain self-similar solutions from a blowup of surface diffusion flow, which one might even expect given initial data with zero enclosed volume, such as a symmetric figure eight.*

We now need to show that \tilde{f} is nontrivial. The arguments below follow [14], although here the area bound does not present any difficulty.

Lemma 26. *The blowup \tilde{f} constructed above is not a union of planes.*

Proof. Due to the smooth convergence in (26) and the second conclusion in Lemma 24 we have

$$\int_{\tilde{f}^{-1}(B_1(0))} |A|^2 d\mu \geq \epsilon_1 > 0.$$

\square \square

Lemma 27. *If the blowup \tilde{f} constructed above contains a compact connected component C , then $\tilde{M} = C$ and M is diffeomorphic to C .*

Proof. For j sufficiently large, $\phi_j(C)$ is open and closed in M . By the connectedness of M we have $M = \phi_j(C)$ and thus $\tilde{M} = C$. \square \square

Theorem 28 (Nontriviality of the blowup). *Suppose $f : M^2 \times [0, T] \rightarrow \mathbb{R}^3$ is a surface diffusion flow satisfying (4) and let \tilde{f} be the blowup constructed above. Then none of the components of \tilde{f} are compact. In particular, the blowup has a component which is a noncompact nonumbilic surface with $\Delta H \equiv 0$.*

Proof. Observe that the surface area $\mu(M)$ is uniformly bounded away from zero, first by noting that each f_t is an embedding (by Proposition 12) and then by combining the isoperimetric inequality with the computation (2).

Assume that there is a compact component of \tilde{f} . Then Lemma 27 implies that $\tilde{f} : \tilde{M} \rightarrow \mathbb{R}^3$ has no further components. Therefore the surface area of the blowup is bounded. The measure behaves under scaling by

$$\mu(M) \Big|_{t=t_j} = r_j^2 \mu_j(M) \Big|_{t=0}$$

and so we have

$$\mu(M)\Big|_{t=T} = \lim_{j \rightarrow \infty} \mu(M)\Big|_{t=t_j} = \lim_{j \rightarrow \infty} r_j^2 \mu_j(M)\Big|_{t=0} = 0.$$

This is in direct contradiction with the fact that area is uniformly bounded away from zero. Thus there are no compact components of \tilde{f} , and Lemma 26 gives that there must exist at least one nonumbilic noncompact component of \tilde{f} . \square \square

29. ASYMPTOTIC BEHAVIOUR

Combining the analysis of the previous sections we can finally rule out concentration of curvature in finite time.

Proposition 30. *Suppose $f : M^2 \times [0, T) \rightarrow \mathbb{R}^3$ is a (1) flow. Then there exists a constant $\epsilon_0 > 0$ such that if*

$$\int_M |A^\circ|^2 d\mu \Big|_{t=0} < \epsilon_0$$

then $T = \infty$.

Proof. Assume otherwise, and then by the Lifespan Theorem there exists a $T < \infty$ such that curvature concentrates at time T . Performing a blowup as in Section 3 at T we recover a stationary surface \tilde{f} with small tracefree curvature due to the scale invariance of $\|A^\circ\|_2^2$ and Proposition 12. Now Theorem 10 implies \tilde{f} must be a plane or sphere, which contradicts the nontriviality of the blow up, Theorem 28. Thus there does not exist a finite time when curvature concentrates, and so $T = \infty$. \square \square

Smooth (sub)convergence to round spheres now follows by using a similar argument, constructively instead of for the purposes of contradiction. This is similar to Lemma 5.4 in [14], and so we omit the proof.

Lemma 31. *Suppose $f : M^2 \times [0, T) \rightarrow \mathbb{R}^3$ is a surface diffusion flow satisfying (4). Then for any sequence $t_j \nearrow \infty$ there exist $x_j \in \mathbb{R}^3$ and $\phi_j \in \text{Diff}(M)$ such that, after passing to a subsequence, the immersions $f(\phi_j, t) - x_j$ converge smoothly to an embedded round sphere.*

The above implies that Proposition 12 may in fact be strengthened to

$$\int_M |A^\circ|^2 d\mu \searrow 0 \text{ as } t \nearrow \infty.$$

We finish by proving exponential decay of curvature.

Proposition 32. *Suppose $f : M^2 \times [0, T) \rightarrow \mathbb{R}^3$ is a surface diffusion flow satisfying (4). Then there exists a $\lambda > 0$ such that as $t \nearrow \infty$ the following asymptotic statements hold:*

$$\|\nabla_{(k)} A\|_\infty \leq c^k e^{-\lambda t}, \quad \text{and} \quad \|A^\circ\|_\infty \leq c^0 e^{-\lambda t},$$

for $k \geq 1$.

Proof. First note that using an argument similar to that used to prove (17) one may establish the estimate

$$(27) \quad \frac{d}{dt} \int_M |A^\circ|^2 d\mu + \frac{1}{100} \int_M (|\nabla_{(2)} A|^2 + |\nabla A|^2 H^2 + |A^\circ|^2 H^4) d\mu \leq 0.$$

Let $\mathcal{A} = \mu(M_t)|_{t=T}$. Then Lemma 31 above implies that the sectional curvature and mean curvature satisfy

$$\|K\|_\infty \rightarrow \frac{4\pi}{\mathcal{A}}, \quad \text{and} \quad \|H^2\|_\infty \rightarrow \frac{16\pi}{\mathcal{A}}$$

as $t \nearrow \infty$ respectively. Therefore there exists a $t_H < \infty$ such that

$$H^2 \geq c_H > 0, \quad \text{for all } t \geq t_H.$$

From now on we assume $t \geq t_H$. Note that we may assume $c_H < 1$. Using (27),

$$\frac{d}{dt} \int_M |A^\circ|^2 d\mu + \frac{c_H^2}{100} \int_M |\nabla_{(2)} A|^2 + |\nabla A|^2 + |A^\circ|^2 d\mu \leq 0.$$

Integrating gives

$$(28) \quad \int_M |A^\circ|^2 d\mu + \frac{c_H^2}{100} \int_t^\infty \int_M |\nabla_{(2)} A|^2 + |\nabla A|^2 d\mu d\tau \leq e^{-2\lambda t},$$

where $\lambda = \frac{c_H^2}{200}$. Note that we needed Lemma 31 for $\int_M |A^\circ|^2 d\mu|_{t=T} = 0$ in the above. From this estimate and again Lemma 31 we can obtain exponential decay. From Proposition 17, taking $\rho \nearrow \infty$ gives

$$\frac{d}{dt} \int_M |\nabla_{(k)} A|^2 d\mu + \int_M |\nabla_{(k+2)} A|^2 d\mu \leq \int_M (P_3^{k+2}(A) + P_5^k(A)) * \nabla_{(k)} A d\mu.$$

Note that the term $c\|A\|_{2, [\gamma>0]}^2$ disappeared due to the dependence of the constant on ρ . From Lemma 31 we know that A and all its derivatives remain bounded as $t \nearrow \infty$, so we estimate

$$\begin{aligned} \int_M P_2^0(A) * \nabla_{(k+2)} A * \nabla_{(k)} A d\mu &\leq \epsilon \int_M |\nabla_{(k+2)} A|^2 d\mu + c_\epsilon \int_M |\nabla_{(k)} A|^2 d\mu, \\ \int_M (\tilde{P}_3^{k+2}(A) + P_5^m(A)) * \nabla_{(k)} A d\mu &\leq c \sum_{j=1}^{k+1} \int_M |\nabla_{(j)} A|^2 d\mu, \end{aligned}$$

where the constant is not universal (i.e. depends on derivatives of curvature). $\tilde{P}_3^{k+2}(A)$ denotes all terms of type $P_3^{k+2}(A)$ that do not contain the $(k+2)$ -th derivative. We thus obtain

$$\frac{d}{dt} \int_M |\nabla_{(k)} A|^2 d\mu + \frac{1}{2} \int_M |\nabla_{(k+2)} A|^2 d\mu \leq c \sum_{j=1}^{k+1} \int_M |\nabla_{(j)} A|^2 d\mu.$$

A proof by induction using (28) then gives

$$\|\nabla_{(k)} A\|_2^2 + \frac{c_H^2}{100} \int_t^\infty \|\nabla_{(k+2)} A\|_2^2 d\tau \leq e^{-2\lambda t}.$$

This gives us the estimates

$$\|A^\circ\|_2 \leq ce^{-\lambda t}, \quad \text{and} \quad \|\nabla_{(k)} A\|_2 \leq ce^{-\lambda t}.$$

Using Lemma 7 finishes the proof. \square \square

REFERENCES

- [1] Baras, P., Duchon, J., Robert, R.: Evolution D'une Interface par diffusion de surface. *Communications in Partial Differential Equations* **9**(4), 313–335 (1984)
- [2] Blatt, S.: Loss of convexity and embeddedness for geometric evolution equations of higher order. *Journal of Evolution Equations* **10**(1), 21–27 (2010)
- [3] Cahn, J., Elliott, C., Novick-Cohen, A.: The Cahn-Hilliard equation with a concentration dependent mobility: motion by minus the Laplacian of the mean curvature. *European J. Appl. Math* **7**(3), 287–301 (1996)
- [4] Cahn, J., Taylor, J.: Surface motion by surface diffusion. *Acta metallurgica et materialia* **42**(4), 1045–1063 (1994)
- [5] Chruściel, P.: Semi-global existence and convergence of solutions of the Robinson-Trautman (2-dimensional Calabi) equation. *Communications in Mathematical Physics* **137**(2), 289–313 (1991)
- [6] Davì, F., Gurtin, M.: On the motion of a phase interface by surface diffusion. *Zeitschrift für Angewandte Mathematik und Physik (ZAMP)* **41**(6), 782–811 (1990)
- [7] Elliott, C., Garcke, H.: Existence results for diffusive surface motion laws. *Adv. Math. Sci. Appl* **7**(1), 465–488 (1997)
- [8] Escher, J., Mayer, U.F., Simonett, G.: The surface diffusion flow for immersed hypersurfaces. *SIAM Journal on Mathematical Analysis* **29**(6), 1419–1433 (1998)
- [9] Fife, P.: Models for phase separation and their mathematics. *Electronic Journal of Differential Equations* **2000**(48), 1–26 (2000)
- [10] Hamilton, R.S.: Three-manifolds with positive Ricci curvature. *J. Differential Geom.* **17**(2), 255–306 (1982)
- [11] Huisken, G.: Flow by mean curvature of convex surfaces into spheres. *J. Differential Geom* **20**(1), 237–266 (1984)
- [12] Huisken, G., Polden, A.: Geometric evolution equations for hypersurfaces. *Calculus of variations and geometric evolution problems (Cetraro, 1996)*, Springer, Berlin pp. 45–84 (1999)
- [13] Ito, K.: The surface diffusion flow equation does not preserve the convexity. *RIMS Kokyuroku* **1105**, 10–21 (1999)
- [14] Kuwert, E., Schatzle, R.: The Willmore flow with small initial energy. *J. Differential Geom* **57**(3), 409–441 (2001)
- [15] Kuwert, E., Schatzle, R.: Gradient flow for the Willmore functional. *Comm. Anal. Geom* **10**(2), 307–339 (2002)
- [16] Li, P., Yau, S.T.: A new conformal invariant and its applications to the Willmore conjecture and the first eigenvalue of compact surfaces. *Invent. Math.* **69**, 269–291 (1982)
- [17] Mantegazza, C.: Smooth geometric evolutions of hypersurfaces. *Geometric And Functional Analysis* **12**(1), 138–182 (2002)
- [18] Mayer, U.: Numerical solutions for the surface diffusion flow in three space dimensions. *Computational and Applied Mathematics* **20**(3), 361–379 (2001)
- [19] Michael, J., Simon, L.: Sobolev and mean-value inequalities on generalized submanifolds of \mathbb{R}^n . *Comm. Pure Appl. Math* **26**(36), 1–379 (1973)
- [20] Mullins, W.: Theory of thermal grooving *J. Appl. Phys* **28**, 333–9 (1957)
- [21] Polden, A.: Curves and surfaces of total curvature and fourth order flows. Ph.D. thesis, Thesis, Mathematisches Institut, Uni. Tübingen (1996)
- [22] Simon, L.: Existence of surfaces minimizing the Willmore functional. *Comm. Anal. Geom* **1**(2), 281–326 (1993)
- [23] Simonett, G.: The Willmore Flow for Near Spheres. *Differential and Integral Equations* **14**(8), 1005–1014 (2001)
- [24] Wheeler, G.: Fourth order geometric evolution equations. Ph.D. thesis, University of Wollongong (2010)
- [25] Wheeler, G., McCoy, J., Williams, G.: Lifespan theorem for constrained surface diffusion flows. to appear in *Math. Z.* (2010)