

University of Wollongong

Research Online

Australian Institute for Innovative Materials -
Papers

Australian Institute for Innovative Materials

2012

Flux pinning and vortex transitions in doped BaFe₂As₂ single crystals

S R. Ghorbani

Hakim Sabzevari University, University Of Wollongong, ghorbani@uow.edu.au

X L. Wang

University of Wollongong, xiaolin@uow.edu.au

M Shabazi

University of Wollongong

S X. Dou

University of Wollongong, shi@uow.edu.au

K Y. Choi

Sogang University, Seoul National University

See next page for additional authors

Follow this and additional works at: <https://ro.uow.edu.au/aiimpapers>



Part of the [Engineering Commons](#), and the [Physical Sciences and Mathematics Commons](#)

Research Online is the open access institutional repository for the University of Wollongong. For further information contact the UOW Library: research-pubs@uow.edu.au

Flux pinning and vortex transitions in doped BaFe₂As₂ single crystals

Abstract

The vortex liquid-to-glass transition has been studied in Ba_{0.72}K_{0.28}Fe₂As₂ (BaK-122), Ba(Fe_{0.91}Co_{0.09})₂As₂ (BaCo-122), and Ba(Fe_{0.95}Ni_{0.05})₂As₂ (BaNi-122) single crystal with superconducting transition temperature, $T_c=31.7, 17.3,$ and 18 K, respectively, by magnetoresistance measurements. For temperatures below T_c , the resistivity curves were measured in magnetic fields within the range of $0 \leq B \leq 13$ T, and the pinning potential was scaled according to a modified model for vortex liquid resistivity. Good scaling of the resistivity $\rho(B, T)$ and the effective pinning energy $U_0(B, T)$ were obtained. The vortex state is three-dimensional at temperatures lower than a characteristic temperature T^* . The vortex phase diagram was determined based on the evolution of the vortex-glass transition temperature T_g with magnetic field and the upper critical field, H_{c2} . We found that non-magnetic K doping results in a high glass line close to the H_{c2} , while magnetic Ni and Co doping causes a low glass line which is far away from the H_{c2} . Our results suggest that non-magnetic induced disorder is more favourable for enhancement of pinning strength compared to magnetic induced disorder. Our results show that the pinning potential is responsible for the difference in the glass states

Keywords

crystals, flux, vortex, pinning, transitions, doped, single, bafe2as2

Disciplines

Engineering | Physical Sciences and Mathematics

Publication Details

Ghorbani, S. R., Wang, X. L., Shabazi, M., Dou, S. X., Choi, K. Y. & Lin, C. T. (2012). Flux pinning and vortex transitions in doped BaFe₂As₂ single crystals. *Applied Physics Letters*, 100 (7), 072603-1-072603-4.

Authors

S R. Ghorbani, X L. Wang, M Shabazi, S X. Dou, K Y. Choi, and C T. Lin

Flux pinning and vortex transitions in doped BaFe₂As₂ single crystals

S. R. Ghorbani, X. L. Wang, M. Shabazi, S. X. Dou, K. Y. Choi et al.

Citation: *Appl. Phys. Lett.* **100**, 072603 (2012); doi: 10.1063/1.3685507

View online: <http://dx.doi.org/10.1063/1.3685507>

View Table of Contents: <http://apl.aip.org/resource/1/APPLAB/v100/i7>

Published by the American Institute of Physics.

Related Articles

Competition of 3d/4f orbitals due to competing conductivity and ferromagnetism in Fe/CoAs layers in Eu(Fe_{0.89}Co_{0.11})₂As₂

J. Appl. Phys. **113**, 013907 (2013)

Enhanced anisotropic heat conduction in multi-walled carbon nanotubes

J. Appl. Phys. **113**, 014301 (2013)

In-plane electronic confinement in superconducting LaAlO₃/SrTiO₃ nanostructures

Appl. Phys. Lett. **101**, 222601 (2012)

Effect of Co-doping on the resistivity and thermopower of SmFe_{1-x}CoxAsO (0.0 ≤ x ≤ 0.3)

AIP Advances **2**, 042137 (2012)

Critical current reduction in coated conductors when in-plane fields are applied

J. Appl. Phys. **112**, 073918 (2012)

Additional information on *Appl. Phys. Lett.*

Journal Homepage: <http://apl.aip.org/>

Journal Information: http://apl.aip.org/about/about_the_journal

Top downloads: http://apl.aip.org/features/most_downloaded

Information for Authors: <http://apl.aip.org/authors>

ADVERTISEMENT

AIP | Applied Physics
Letters

**EXPLORE WHAT'S
NEW IN APL**

SUBMIT YOUR PAPER NOW!

SURFACES AND INTERFACES
Focusing on physical, chemical, biological, structural, optical, magnetic and electrical properties of surfaces and interfaces, and more...

ENERGY CONVERSION AND STORAGE
Focusing on all aspects of static and dynamic energy conversion, energy storage, photovoltaics, solar fuels, batteries, capacitors, thermoelectrics, and more...

Flux pinning and vortex transitions in doped BaFe₂As₂ single crystals

S. R. Ghorbani,^{1,2} X. L. Wang,^{1,a)} M. Shabazi,¹ S. X. Dou,¹ K. Y. Choi,³ and C. T. Lin⁴

¹*Institute for Superconducting and Electronic Materials, University of Wollongong, Wollongong, New South Wales 2522, Australia*

²*Department of Physics, Sabzevar Tarbiat Moallem University, P.O. Box 397, Sabzevar, Iran*

³*Frontier Physics Research Division and Department of Physics and Astronomy, Seoul National University, Seoul 151-747, Korea*

⁴*Max Planck Institute for Solid State Research, Heisenbergstr 1, 70569 Stuttgart, Germany*

(Received 30 October 2011; accepted 30 January 2012; published online 17 February 2012)

The vortex liquid-to-glass transition has been studied in Ba_{0.72}K_{0.28}Fe₂As₂ (BaK-122), Ba(Fe_{0.91}Co_{0.09})₂As₂ (BaCo-122), and Ba(Fe_{0.95}Ni_{0.05})₂As₂ (BaNi-122) single crystal with superconducting transition temperature, $T_c = 31.7$, 17.3 , and 18 K, respectively, by magnetoresistance measurements. For temperatures below T_c , the resistivity curves were measured in magnetic fields within the range of $0 \leq B \leq 13$ T, and the pinning potential was scaled according to a modified model for vortex liquid resistivity. Good scaling of the resistivity $\rho(B, T)$ and the effective pinning energy $U_0(B, T)$ were obtained. The vortex state is three-dimensional at temperatures lower than a characteristic temperature T^* . The vortex phase diagram was determined based on the evolution of the vortex-glass transition temperature T_g with magnetic field and the upper critical field, H_{c2} . We found that non-magnetic K doping results in a high glass line close to the H_{c2} , while magnetic Ni and Co doping causes a low glass line which is far away from the H_{c2} . Our results suggest that non-magnetic induced disorder is more favourable for enhancement of pinning strength compared to magnetic induced disorder. Our results show that the pinning potential is responsible for the difference in the glass states. © 2012 American Institute of Physics. [doi:10.1063/1.3685507]

In the mixed states of type II superconductors, vortices form the Abrikosov lattice as a result of quantization of magnetic flux. Depending on the magnetic field, the vortex lattice can change from the solid to the glass state and will further transform to a liquid state and disappear at $H > H_{c2}$, where H_{c2} is the upper critical field. The boundary between the lattice and glass or the glass and liquid phases is strongly affected by both the anisotropy and the disorder of a superconductor. The vortex glass line can be significantly suppressed in highly anisotropic superconductors such as high superconducting transition temperature (T_c) cuprates, as the field required for the glass state is inversely proportional to the square of the anisotropy.¹ Local disorder is another important factor affecting the vortex glass line. In high- T_c cuprates, vortex lattices are soft and form two-dimensional pancake vortices, which are easily melted into vortex liquid by magnetic field and thermal fluctuations. Therefore, the glass line is far below H_{c2} in cuprates. From the viewpoint of applications, the glass line determines the critical field below which the vortices can be pinned or supercurrent can still survive. The Fe-based superconductors are found to show a small anisotropy in REFeAsO_{1-x}F_x (RE-1111 phase, with RE a rare earth element)²⁻⁵ and doped BaFe₂As₂ (122 phase).^{6,7} It has been found that doping the latter with K, a non-magnetic dopant occupying Ba sites, results in a maximum T_c of 35 K and induces a very strong intrinsic pinning strength with high critical current density (J_c) and H_{c2} simultaneously.⁷ Whereas, magnetic dopants, such as Co and Ni, lead to a low T_c of about 18–22 K and lower J_c and H_{c2} than

for K doped 122, although the grain boundaries are not detrimental to the J_c , as has been reported for Co-doped 122 epitaxial thin films.⁸ So far, little work has been reported on the determination of glass lines in the pnictide compounds. These facts have motivated us to raise a fundamental question as to whether or not disorder induced by magnetic or non-magnetic dopants can cause great differences in the glass lines, since the anisotropy is only around 2–3 in 122 compounds. In this work, the vortex liquid-to-glass transition has been studied in Ba_{0.72}K_{0.28}Fe₂As₂ (BaK-122), Ba(Fe_{0.91}Co_{0.09})₂As₂ (BaCo-122), and Ba(Fe_{0.95}Ni_{0.05})₂As₂ (BaNi-122) single crystal with $T_c = 31.7$, 17.3 , and 18 K, respectively, by magnetoresistance measurements. We found that non-magnetic K doping results in a high glass line close to the H_{c2} , while magnetic Ni and Co doping causes a low glass line which is far away from the H_{c2} . Our results suggest that non-magnetic induced disorder is more favourable for enhancement of pinning strength compared to magnetic induced disorder.

The phase transition can be clearly visualized through the difference in resistive behaviour between the solid and liquid phases. The broadening of the resistivity transition in magnetic field is a direct consequence of the thermal fluctuation in the vortex system. Therefore, resistive transport measurements are commonly used to study vortices and vortex phase transitions.^{1,9-13} The vortex phase transition in cuprate superconductors can be understood using vortex-glass theory. According to this theory,⁹ in the vortex glass state and close to the glass transition temperature T_g , the resistivity decreases as a power law

$$\rho = \rho_0 \left| \frac{T}{T_g} - 1 \right|^s, \quad (1)$$

^{a)} Author to whom correspondence should be addressed. Electronic mail: xiaolin@uow.edu.au.

where s is a constant related to the various types of disorder, which can be introduced by columnar defects, point defects, dopants, and vacancies. ρ_0 is the characteristic resistivity of the normal state. Equation (1) can be modified by using the energy difference $k_B T - U_0$, where U_0 is the effective pinning energy. Therefore, the driving force of the transition depends on the energy difference, instead of the temperature difference of $T - T_g$.¹¹ Therefore, Eq. (1) is re-formulated as

$$\rho = \rho_n \left| \frac{k_B T}{U_0(B, T)} - 1 \right|^s. \quad (2)$$

Here ρ_n is the normal state resistivity at the onset of the transition. In this model, the vortex solid to vortex liquid transition occurs when the two energy scales are equal, $U_0(B, T_g) = k_B T_g$. An empirical effective pinning energy was found¹⁴

$$U_0(B, T) = U_B \left(1 - \frac{T}{T_c} \right) \quad \text{with} \quad U_B = k_B T_c / (B/B_0)^\beta, \quad (3)$$

where both B_0 , which is inversely proportional to the square of the mass anisotropy, and β are temperature and field-independent constants. By considering the effective pinning energy at the glass temperature, i.e., $U_0(B, T_g) = k_B T_g$, the temperature dependent vortex glass line is obtained from^{10,11}

$$B_g(T) = B_0 \left(\frac{1 - T/T_c}{T/T_c} \right)^{1/\beta}. \quad (4)$$

This has been used to determine the vortex glass for Y-123.¹⁰⁻¹³

A useful scaling form for resistivity was obtained by combination of Eqs. (2) and (3)

$$\rho = \rho_n \left| \frac{T(T_c - T_g)}{T_g(T_c - T)} - 1 \right|^s. \quad (5)$$

Therefore, the resistivity in the vortex glass state depends on magnetic field through T_g . Although Strachan *et al.*¹⁵ suggested that agreement with scaling alone does not prove the existence of a phase transition, but experimental evidence for the vortex-glass phase has been obtained for various high- T_c superconductors such as Y-123,¹⁰⁻¹² Tl-2212,^{13,14} MgB₂,¹⁶ and Ba_{0.55}K_{0.45}Fe₂As₂ (Ref. 17) through the scaling exponent, as predicted by vortex-glass and modified vortex-glass theory. However, the vortex-glass phase transition and scaling behaviour have not been investigated systematically in detail in the Fe-based superconductors.

Here, we have studied the glass transition from the vortex liquid side in BaK-122, BaCo-122, and BaNi-122 single crystal by magnetoresistance measurements. Good scaling of the $\rho(B, T)$ and $U_0(B, T)$ was obtained with the critical exponents s for all crystals. The vortex phase diagram was determined based on the evolution of the vortex-glass transition temperature T_g with magnetic field and the upper critical field. We have also compared the modified vortex-glass model and the thermally activated flux flow model in the vortex liquid region.

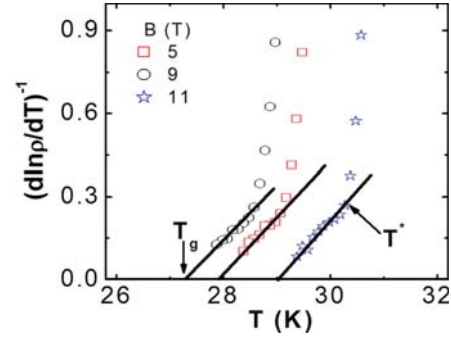


FIG. 1. (Color online) Determination of the glass transition temperature from the Vogel-Fulcher relation for BaK-122 at three different magnetic fields.

Single crystals with the nominal composition BaK-122, BaCo-122, and BaNi-122 were prepared by a self-flux method. Details of the single crystal growth are described elsewhere.¹⁸ The as-grown single crystals were cleaved and shaped into rectangular bars for measurements. The transport properties were measured over a wide range of temperature and magnetic fields up to 13 T using Physical Properties Measurement System (PPMS, Quantum Design).

According to the vortex glass model, Eq. (1), the resistance goes to zero at the glass temperature T_g as $\rho \propto (T - T_g)^s$. Consequently $T_g(B)$ can be extracted by applying the Vogel-Fulcher relation, $(d \ln \rho / dT)^{-1} \propto (T - T_g)$, to the resistive tails. According to Eq. (5), the inverse of the logarithmic derivative is $(\partial \ln \rho / \partial T)^{-1} = (T - T_g / s)(T_c - T / T_c - T_g)$. This expression only differs from the Vogel-Fulcher relation by a correction factor $A(T) = (T_c - T) / (T_c - T_g)$ which is close to one at temperatures sufficiently close to T_g . Therefore, the usual Vogel-Fulcher relation can be applied to estimate T_g directly. As shown in Fig. 1, the low temperature data are consistent with the Vogel-Fulcher relation up to the temperature of T^* , which shows a deviation from the straight line, with an intercept of T_g for BaK-122, BaCo-122, and BaNi-122. At temperature T^* , vortices change from 2D to 3D. In layered superconductors above this temperature, the vortex flux lines act as a two-dimensional, 2D, pancake vortex.¹⁹

The obtained s , the inverse of the slope of the straight line, is 6.0 ± 0.5 , 5.3 ± 0.3 , and 5.1 ± 0.4 for BaK-122, BaCo-122, and BaNi-122 single crystals, respectively. For BaK-122, the s parameter is roughly three times larger than that obtained for Ba_{0.55}K_{0.45}Fe₂As₂.¹⁷

The electrical resistance of BaK-122, BaCo-122, and BaNi-122 single crystals, measured at fields up to 13 T for

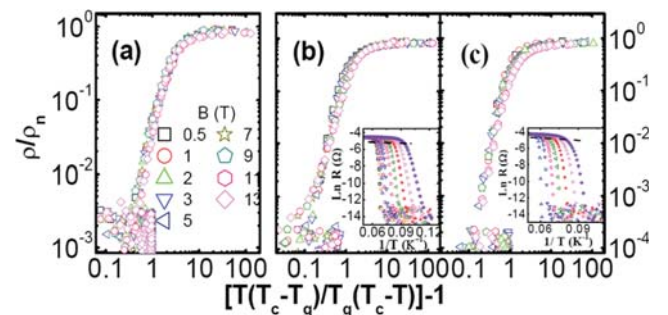


FIG. 2. (Color online) Resistivity scaling according to Eq. (5) for (a) BaK-122, (b) BaCo-122, and (c) BaNi-122 single crystal for $0 \leq B \leq 13$ T. Inset: Arrhenius plot of the resistance.

$H//c$, is shown in the Arrhenius plot in inset Fig. 2. The transition temperatures, T_c , are 31.7, 17.3, and 18.0 K, with almost the same small transition width (ΔT_c) of 0.7 ± 0.1 K for BaK-122, BaCo-122, and BaNi-122 single crystals, respectively.

Equation (5) suggests a scaling behaviour of the vortex glass between the normalized resistivity ρ/ρ_n and the scaled temperature $T_{sc} = [T(T_c - T_g)/T_g(T_c - T)] - 1$. Figure 2 shows the ρ/ρ_n versus scaling temperature T_{sc} . As can be seen in Fig. 2, the resistivity transition at fields between 0 and 13 T has been scaled into one curve with the obtained critical exponent Vortex glass with the same resistivity scaling behaviour has been observed in several high- T_c superconductors such as oxygen-deficient Y-123 single crystal^{11,12} and Tl-2212 thin film.¹³

The upper critical field, B_{c2} , is obtained from the 90% values of its corresponding resistivity transition. Using the estimated vortex glass line B_g , the B - T phase diagram for the BaK-122, BaCo-122, and BaNi-122 single crystals is shown in Fig. 3. According to the collective pinning model,²⁰ the disorder-induced spatial fluctuations in the solid-vortex lattice can be clearly divided into markedly different regimes according to the strength of the applied field. As is shown in the inset of Fig. 3, two different regimes are distinguishable: (1) vortex glass, which governs the region below B_g ; (2) vortex liquid, which holds between B_g and B_{c2} , where thermal fluctuations are important. As can be seen from Fig. 3, the vortex-glass phase indicates that the K doped BaFe₂As₂ single crystal has a very narrow region of the vortex-liquid phase, which is denoted by ΔT in the inset of Fig. 3, with a small ΔT of 0.06 ± 0.02 T and at magnetic field of 0.5 up to 13 T (weakly field dependent), which has originated from the vastly enhanced vortex pinning. However, the glass lines for both Co ($\Delta T = 0.18 \pm 0.02$) and Ni doped 122 single crystals are far from the H_{c2} lines and strongly temperature dependent for Ni122 ($\Delta T = 0.03$ – 0.22 for the $B = 0.5$ – 13 T range). This indicates that the vortex liquid state exists in a wide range of fields below H_{c2} . This is significantly different from BaK-122. The Co and Ni (magnetic doping) shifts the vortex-glass transition to lower temperature or softening vortex lattice as compared to K (non-magnetic doping). The magnetic doping was decreased T_c and the slope of the

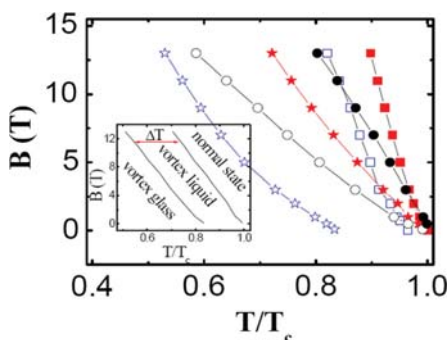


FIG. 3. (Color online) Phase diagrams of BaK-122 (squares), BaCo-122 (stars), and BaNi-122 (circles) single crystals. B_g was obtained from the experimental $\rho(T)$ and $U_0(T)$ data (see Figs. 2 and 4). $B_{c2}(T)$ values were obtained from the criterion of the 90% values of their corresponding resistivity transition. Open symbols: B_{c2} and solid symbols: B_g . Inset: general phase diagram.

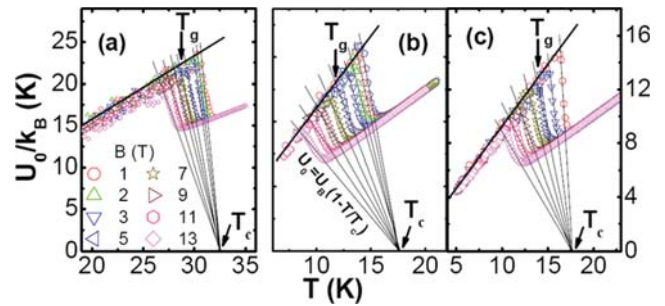


FIG. 4. (Color online) The pinning potential $U_0(B,T)$ as calculated according to Eq. (6) in field of $0 \leq H \leq 13$ T for (a) BaK-122, (b) BaCo-122, and (c) BaNi-122 single crystal.

mean field normal phase-mixed phase boundary line, $dH_{c2}/dT|_{T=T_c}$, as seen in Figs. 2 and 3. The smaller change of coherence length ($\xi = (\phi_0/2\pi H_{c2})^{1/2}$, ϕ_0 is the flux quanta) by non-magnetic doping relative to magnetic doping implies some possible vortex-impurity interactions causing an enhanced pinning or stiffening of the vortex system and bring about the transition to the vortex-solid phase at higher temperature. This is supported by higher pinning potential for K doping than that for Co and Ni as will be discussed later. As Ni and Co substitute Fe sites, it can induce distortion or imperfection to the FeAs planes compared to the K doping which can mainly provide more charge carriers, but cause little changes to the FeAs planes. It seems that the distortion of FeAs plane is likely associated with pinning and T_c . A further study is needed to elucidate the origin of the differences of the pinning in non-magnetic K and magnetic Ni and Co ion doped 122 crystals. Another important point is that one can directly observe that ΔT of BaK-122 is smaller than for BaCo-122 or BaNi-122, which is an indication that the pinning potential of BaK-122 is stronger than for either BaCo-122 or BaNi-122 single crystals. This is supported by the pinning potential results, as can be seen in Fig. 5. The results show that non-magnetic doping, such as K on Ba site, and magnetic doping, such as Co or Ni on the Fe site, have different effects on the vortex glass line. The question is what the effect will be on the vortex glass line of magnetic and non-magnetic substitution on the same site, especially the Fe site, which is located on the FeAs plane. This requires further study.

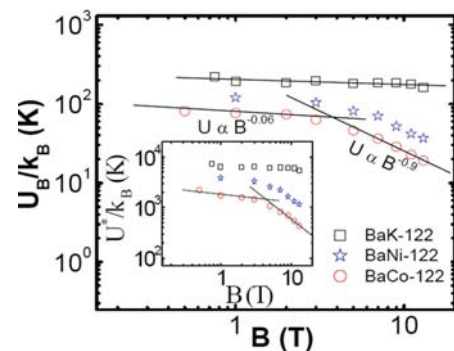


FIG. 5. (Color online) Magnetic field dependence of U_B/k_B as estimated from Eq. (3) and the slopes in Fig. 4. The solid lines are a fit to the data giving the relation $U_B/k_B \approx B^{-n}$ with $n = 0.06$ at $B < 3$ T and $n = 0.9$ at $B > 3$ T. Inset: Field dependence of the estimated activation energy based on the TAFF model.

By solving Eq. (2) for U_0 , one obtains

$$U_0(B, T) = k_B T \left[1 + \left(\frac{\rho}{\rho_n} \right)^{1/s} \right]^{-1}. \quad (6)$$

According to Eq. (6), the pinning potential can be calculated directly from experimental data, provided that one knows the ρ_n and the exponent s . The s parameter is estimated as the inverse slope of the resistivity in the vortex glass state, as mentioned above, and the ρ_n has been taken as the normal state resistivity at $T=35$ K for BaK-122 and $T=20$ K for both BaCo-122 and BaNi-122 single crystals. The calculated $U_0(B, T)$ of the BaK-122, BaCo-122, and BaNi-122 single crystals for fields between 0 and 13 T is shown in Fig. 4. As clearly shown by the solid lines, the low resistivity of the curves is well described by Eq. (3) with a field dependent U_B .

As can be seen in Fig. 4, an extrapolation of the linear behaviour of $U_0(T)$ at temperatures close to T_g and at different fields merges in the point $T=T_c$ and $U_0=0$ which is in good agreement with Eq. (3). Therefore, the field dependence of the pinning energy U_B can be found directly from the slope of these lines. The magnetic dependence of the estimated result from Fig. 4 for U_B is shown in Fig. 5. As can be seen in Fig. 5, the pinning potential of BaK-122 single crystal is constant, while different power law dependence is found at low and high magnetic fields for both BaCo-122, and BaNi-122 single crystal. The U_B/k_B decreases slowly with increasing applied magnetic field for $B < 3$ T, scaled as $B^{-0.06}$, and then decreases as $B^{-0.9}$ for $B > 3$ T. This result suggests that the single vortex pinning may co-exist with collective creep in low fields and then the collective creep dominates in high magnetic fields.

Another important point is that one can directly obtain T_g by considering the crossing points of the lines $U_0(B, T)$ and the line $U_0 = k_B T$, as indicated by the arrows in Fig. 4. This criterion was used for estimating T_g in the modified vortex-glass model.^{11,12} Therefore, $U_0(B, T)$ is the average pinning energy in the system and is responsible for the vortex solid to liquid transition when $U_0(B, T_g) = k_B T_g$, as discussed in the Introduction.

In many studies, the vortex liquid resistivity is described by a thermally activated flux flow (TAFF) model, $\rho(T, H) = \rho_n \exp(U^*/k_B T)$,²¹ where ρ_n is the normal state resistivity and k_B is the Boltzmann constant. By using the TAFF model, we obtained the activation energy, $U^*/k_B = \partial \ln \rho / \partial (1/T)$, which is shown in the inset of Fig. 5. At $B = 1$ T, U^*/k_B was 6400, 3311, and 1626 K for the BaK-122, BaCo-122, and BaNi-122 single crystals, respectively, while U_B/k_B was 192, 77, and 32 K. Although U^* is larger than U_B in value, it has a similar magnetic field dependence. For comparison of these values, by taking the derivative $\partial \ln \rho / \partial (1/T)$ of Eq. (2) and using Eq. (3), one can obtain the relation $U^* = \frac{s[1+(\rho/\rho_n)^{1/s}]^2}{(\rho/\rho_n)^{1/s}} U_B = A U_B$. By using s and experimental data of $\rho/\rho_n \approx 10^{-1} - 10^{-3}$, the correction factor A is between 24–33, 21–32, and 20–28 for the BaK-122, BaCo-122, and

BaNi-122 single crystals, respectively. Therefore, there is an excellent agreement between both models.

In conclusion, it was shown that the glass transition introduced based on a modified model for the vortex-glass transition can be applied to the doped BaFe₂As₂ superconductors. For temperature below the superconducting transition temperature, a scaling of all measured resistivity $\rho(B, T)$ and pinning potential $U_0(B, T)$ values in magnetic fields up to 13 T with the critical exponents s is obtained. The vortex phase diagram was determined based on the evolution of the vortex-glass transition temperature T_g with magnetic field and the upper critical field. The results suggest that the vortex region of BaK-122 is shorter than it is for BaCo-122 and BaNi-122, and it is also magnetic field independent in magnetic fields smaller than 13 T. It was found that the pinning force of BaK-122 is stronger than for BaCo-122 and BaNi-122. Furthermore, we compared these results with the thermally activated flux flow behaviour which is usually employed to account for the resistivity in the vortex liquid region.

This work was supported by the Australian Research Council through Discovery Projects DP1094073 and DP0770205.

- ¹J. Deak, M. McElfresh, D. W. Face, and W. L. Holstein, *Phys. Rev. B* **52**, R3880 (1995).
- ²Y. Kamihara, T. Watanabe, M. Hirano, and H. Hosono, *J. Am. Chem. Soc.* **130**, 3296 (2008).
- ³X. H. Chen, T. Wu, G. Wu, R. H. Liu, H. Chen, and D. F. Fang, *Nature (London)* **453**, 761 (2008).
- ⁴Z. A. Ren, J. Yang, W. Lu, W. Yi, X.-L. Shen, Z.-C. Li, G.-C. Che, X.-L. Dong, L.-L. Sun, F. Zhou *et al.*, *EPL* **82**, 57002 (2008).
- ⁵M. Rotter, M. Tegel, and D. Johrendt, *Phys. Rev. Lett.* **101**, 107006 (2008).
- ⁶H. Q. Yuan, J. Singleton, F. F. Balakirev, S. A. Baily, G. F. Chen, J. L. Luo, and N. L. Wang, *Nature (London)* **457**, 565 (2009).
- ⁷X. L. Wang, S. R. Ghorbani, S.-I. Lee, S. X. Dou, C. T. Lin, T. H. Johansen, K.-H. Müller, Z. X. Cheng, G. Peleckis, M. Shabazi *et al.*, *Phys. Rev. B* **82**, 024525 (2010).
- ⁸T. Katase, Y. Ishimaru, A. Tsukamoto, H. Hiramatsu, T. Kamiya, K. Tanabe, and H. Hosono, *Nature Commun.* **2**, 409 (2011).
- ⁹D. S. Fisher, M. P. A. Fisher, and D. A. Huse, *Phys. Rev. B* **43**, 130 (1991).
- ¹⁰B. Lundqvist, A. Rydh, Yu. Eltsev, Ö. Rapp, and M. Andersson, *Phys. Rev. B* **57**, R14064 (1998).
- ¹¹A. Rydh, Ö. Rapp, and M. Andersson, *Phys. Rev. Lett.* **83**, 1850 (1999).
- ¹²M. Andersson, A. Rydh, and Ö. Rapp, *Phys. Rev. B* **63**, 184511 (2001).
- ¹³B. Espinosa-Arronte and M. Andersson, *Phys. Rev. B* **71**, 024507 (2005).
- ¹⁴H. Wen, H. A. Radovan, F. M. Kamm, P. Ziemann, S. L. Yan, L. Fang, and M. S. Si, *Phys. Rev. Lett.* **80**, 3859 (1998).
- ¹⁵D. R. Strachan, M. C. Sullivan, P. Fournier, S. P. Pai, T. Venkatesan, and C. J. Lobb, *Phys. Rev. Lett.* **87**, 067007 (2001).
- ¹⁶H. J. Kim, W. N. Kang, E. I. Choi, M. S. Kim, K. H. P. Kim, and S.-I. Lee, *Phys. Rev. Lett.* **87**, 087002 (2001).
- ¹⁷H. J. Kim, Y. Liu, Y. S. Oh, S. Khim, I. Kim, G. R. Stewart, and K. H. Kim, *Phys. Rev. B* **79**, 014514 (2009).
- ¹⁸A. S. Sefat, R. Jin, M. A. McGuire, B. C. Sales, D. J. Singh, and D. Mandrus, *Phys. Rev. Lett.* **101**, 117004 (2008).
- ¹⁹H. Safar, P. L. Gammel, D. J. Bishop, D. B. Mitzi, and A. Kapitulnik, *Phys. Rev. Lett.* **68**, 2672 (1992).
- ²⁰G. Blatter, M. V. Feigel'man, V. B. Geshkenbein, A. I. Larkin, and V. M. Vinokur, *Rev. Mod. Phys.* **66**, 1125 (1994).
- ²¹P. H. Kes, J. Aarts, J. van den Berg, C. J. van der Beek, and J. A. Mydosh, *Supercond. Sci. Technol.* **1**, 242 (1989).