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Inhomogeneities in YBa₂Cu₃O₇ thin films with reduced thickness

O V. Shcherbakova

University of Wollongong, olga@uow.edu.au

A V. Pan

University of Wollongong, pan@uow.edu.au

S K. Gorman

University Of Wollongong

S A. Fedoseev

University of Wollongong, sergey@uow.edu.au

I A. Golovchanskiy

University of Wollongong, ig684@uowmail.edu.au

See next page for additional authors

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Inhomogeneities in YBa₂Cu₃O₇ thin films with reduced thickness

Abstract

Morphology and physical properties of mono- (YBa₂Cu₃O₇) and multilayered (YBa₂Cu₃O₇/SmBa₂Cu₃O₇/YBa₂Cu₃O₇) superconducting thin films with thickness ranging from 90 nm to 28 nm have been investigated. For both types of samples, the superconducting properties degraded with reduction of film thickness. Structural and electromagnetic properties were visualized through scanning electron microscopy and magneto-optical imaging, respectively, and revealed high level of inhomogeneity for thinner (<58>nm) samples. However, samples with thickness above 58 nm showed enhanced homogeneity, which explains better superconducting characteristics observed in these films. Results of this work demonstrate that multilayering approach performed by conventional laser ablation results in degradation of superconducting properties in films with thickness below 90 nm, although has positive impact on morphology of these films, which is crucial for device fabrication process.

Keywords

7, 3o, 2cu, thickness, yba, reduced, inhomogeneities, films, thin

Disciplines

Life Sciences | Physical Sciences and Mathematics | Social and Behavioral Sciences

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Authors

O V. Shcherbakova, A V. Pan, S K. Gorman, S A. Fedoseev, I A. Golovchanskiy, and S. X. Dou

Inhomogeneities in $\text{YBa}_2\text{Cu}_3\text{O}_7$ thin films with reduced thickness

O. V. Shcherbakova, A. V. Pan, S. K. Gorman, S. A. Fedoseev, I. A. Golovchanskiy and S. X. Dou

Institute for Superconducting and Electronic Materials, University of Wollongong, Northfields Avenue, Wollongong, NSW 2522, Australia

Abstract

Morphology and physical properties of mono- ($\text{YBa}_2\text{Cu}_3\text{O}_7$) and multilayered ($\text{YBa}_2\text{Cu}_3\text{O}_7/\text{SmBa}_2\text{Cu}_3\text{O}_7/\text{YBa}_2\text{Cu}_3\text{O}_7$) superconducting thin films with thickness ranging from 90 nm to 28 nm have been investigated. For both types of samples, the superconducting properties degraded with reduction of film thickness. Structural and electromagnetic properties were visualised through scanning electron microscopy and magneto-optical imaging, respectively, and revealed high level of inhomogeneity for thinner (<60 nm) samples. However, samples with thickness above 60 nm showed enhanced homogeneity, which explains better superconducting characteristics observed in these films. Results of this work demonstrate that multilayering approach performed by conventional laser ablation results in degradation of superconducting properties in films with thickness below 90 nm, although has positive impact on morphology of these films, which is crucial for device fabrication process.

Key words: YBCO thin film, multilayering approach, microstructure, magneto-optical imaging, electromagnetic properties
PACS: 74.70.-w, 74.72.-h, 74.25.Ha

Recently, the ultimate sensitivity for low light level systems has been provided by a new class of nanowire Superconducting Single Photon Detector (SSPD) [1,2]. These devices outperform other superconducting photon counters and are significantly better than commercially available semiconducting Si and InGaAs avalanche photodiodes [3]. To date, single photon sensitive detectors have been based on Nb and NbN superconducting thin films operating at 2.6 K or 4 K [4,5]. There would be operational advantages in utilizing $\text{YBa}_2\text{Cu}_3\text{O}_7$ (YBCO) superconducting thin films operating up to 77 K, which would allow substantially relaxed (by modern cryocoolers) cooling requirements. However, for the potential to be realized, there is a need to develop methods of preparing nano-scale YBCO structures with a high degree of structural perfection.

YBCO films of various thicknesses have been previously extensively investigated [6–9]. Superconductivity in the ultrathin films has shown to be present

for thicknesses as small as 1 unit cell [7], but with the critical parameters which are much less than those on thicker (up to micrometer size) films [9]. However, several preceding studies demonstrated that multilayering approach (i.e alternation of YBCO layers with other superconducting ($\text{NdBa}_2\text{Cu}_3\text{O}_7$ [10–12]) or non-superconducting ($\text{PrBa}_2\text{Cu}_3\text{O}_7$ [7], Nd_2CuO_4 [9], CeO_2 [13]) materials drastically improve superconducting properties and morphology of final YBCO film. In spite of relatively thick (up to 1 μm) samples investigated, this approach may be viable for improving superconducting properties of YBCO films at smaller (a few nm) thickness [14].

Being attracted by the idea of possible fabrication of SSPD from YBCO films [15,16], we have initiated study of structural and electromagnetic functionalities of laser ablated YBCO thin films with reduced (from optimal 250-300 nm for our technological process down to 28 nm) thickness. We have also investigated effect of multilayering approach on properties

of YBCO films at the thickness down to 44 nm.

The epitaxial, *c*-axis oriented monolayer YBCO and multilayered YBCO/SmBCO/YBCO (where SmBCO denotes $\text{SmBa}_2\text{Cu}_3\text{O}_7$) thin films were made by pulsed laser deposition using a 248 nm wavelength, KrF eximer laser on (001) SrTiO_3 (STO) substrate. A substrate temperature of 780°C and O_2 pressure of 300 mbar were kept during deposition. Bulk superconducting targets (YBCO or SmBCO) were ablated for different lengths of time in order to fabricate films with different thicknesses (Table 1). The frequency at which the laser was incident on the targets was 1 Hz for all samples. After deposition, films were annealed at 400°C for 1 h at 1 atm of oxygen. The multilayered films were deposited with equal thickness layers of YBCO, SmBCO, and YBCO in this particular order.

Composition	Deposition time, s	Thickness, nm	Abbreviation
YBCO	500	90 ± 14	Y90
YBCO	350	60 ± 13	Y60
YBCO	250	28 ± 8	Y28
YBCO/SmBCO/YBCO	500	90 ± 15	Y/Sm90
YBCO/SmBCO/YBCO	350	58 ± 7	Y/Sm58
YBCO/SmBCO/YBCO	250	44 ± 7	Y/Sm44

Table 1
Samples investigated in this work.

Superconducting properties (critical temperature, magnetic moment) of the films were investigated in the Magnetic Property Measurement System (MPMS, Quantum Design). Distribution of magnetic flux in the films was studied by Magneto-Optical Imaging (MOI) at a temperature of 10 K. Images were acquired by a computer-controlled CCD camera with the magnetic field applied perpendicular to the sample surface. The morphology of the films was studied with the help of the Scanning Electron Microscope (SEM). The thickness measurements were carried out using a Dektak 6M Stylus Profiler. The errors given in Table 1 are the 95% confidence limits (determined by twice the standard deviation) of the 5 measurements taken for each sample.

The surface morphology of the films studied is presented in Fig. 1. As can be seen, the Y28 sample has large cavities on the surface of the film (shown by arrows in Fig. 1 a). This implies that the amount of ablated particles was insufficient to achieve full covering of the substrate area ($5 \times 5 \text{ mm}^2$). As the

deposition time (and thus the thickness of the film) increases, more material arrives; thus, the film continuously covers the surface of the substrate, and becomes notably smoother (Fig. 1 b, c).

The Y/Sm44 sample shows the most inhomogeneous morphology among all samples studied (Fig. 1 d). Note, the thickness of each layer in this sample is as thin as ~ 14 nm. We speculate that the first YBCO layer at this thickness consists of many separated islands. This growth mode likely affects growth of the following SmBCO and YBCO layers, and results in observed high level of inhomogeneity and roughness of the final film. However, as the thickness of each individual layer increases, the morphology of the films improves (Fig. 1 e, f).

There are no large voids on the surface of ≤ 90 nm thick films in contrast to thicker (≥ 250 nm) YBCO films [10]. Indeed, initially strong adatom-substrate bonding results in layer-by-layer growth of the YBCO film on STO substrate up to the thickness of a few unit cells (~ 6 nm), as confirmed by the Reflection High Energy Electron Diffraction studies [6]. With increasing film thickness up to ~ 9 -19 nm, adatom-adatom bonding prevails leading to the switch from two-dimensional (or layer-by-layer) to three-dimensional (or island) growth occurs [8]. Thus, all samples studied (having thickness in the range 28 - 90 nm) are grown in a three-dimensional mode, i.e. with increasing the film thickness YBCO islands become large in size and coalesce. The latter results in appearance of voids in YBCO samples. Note, no voids formed in Y28 sample. They are started to be seen at the thickness of 60 nm and have the size of about 10 nm (shown by arrows in Fig. 1 b). When the thickness of YBCO film reaches 90 nm, the voids have size in the range from 10 nm to 350 nm (Fig. 1 c).

In contrast, there were no voids obtained in the multilayered films at any thicknesses (Fig. 1 d-f). It is likely due to interrupted growth of island's height caused by multilayering. The improved surface morphology of Y/Sm58 and Y/Sm90 films is well suited for structuring a continuous meander line required for fabrication of SSPD device.

Fig. 2 (a) shows magnetic moment of the films as a function of temperature, from which the onset critical temperature values (T_c) are determined and plotted against the film thickness in Fig. 2 (b). Thinner samples (Y28 and Y/Sm44) have broad T_c transition, which is consistent with highly inhomogeneous structure obtained in these films (Fig. 1 a, b), and likely poor crystallinity of YBCO layers in

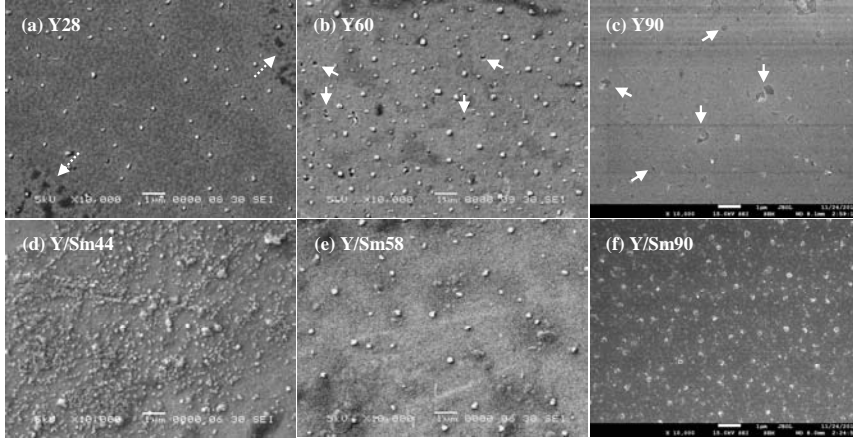


Fig. 1. Morphology of the thin films studied. Images (a, b, c) and (d, e, f) correspond to the different thickness of monolayer and multilayered films, respectively. Arrows in (b) and (c) show some voids formed in 60 nm and 90 nm thick YBCO structures.

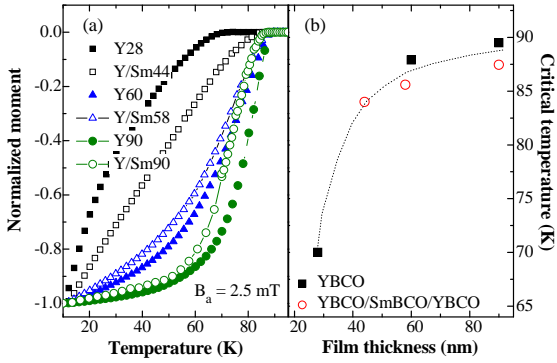


Fig. 2. (a) Normalized magnetic moment as a function of temperature for the samples studied. (b) Critical temperature as a function of film thickness. (The line is a guide for eyes only.)

these thin samples. As the thickness of the films increases, the width of transition becomes narrow, and T_c values increase indicating improved homogeneity in structures formed in these thicker (≥ 60 nm) mono- and multilayered samples (Fig. 1 b, c, e, f).

Fig. 2 (b) demonstrates that there is a gradual reduction of T_c values for films with the thickness varying from 90 nm to 44 nm. After this point (44 nm) there is a dramatic decrease in T_c for the Y28 sample (down to 70 K). Although these results are limited in quantity, they show as the size of the films decreases there is a definite degradation of the superconducting properties, which is more pronounced for thicknesses below approximately 44 nm.

Direct analysis of T_c values suggests that multilayering causes reduction of the eventual supercon-

ducting properties for films with thickness less than 90 nm. However, this may not be the case. Indeed, each single layer building the multilayered structure has the thickness of ≤ 30 nm, which is comparable to thickness of Y28 sample. Thus, each of YBCO layers should have about the same level of structural inhomogeneity (Fig. 1 a) and have broad width of T_c transition as observed for sample Y28 (Fig. 2 a). In reality, the structural and superconducting properties of multilayered samples are gradually enhanced. We speculate that if we are able to find the way to improve crystallinity and continuity of the initial YBCO layer, we may benefit from multilayering approach for thinner (< 90 nm) samples in a similar way as we achieved for thicker films [10,12].

Fig. 3 shows the zero-field magnetization per sample volume (M) for each film at 10 K, 50 K and 77 K. From this graph the trend in critical current density can be determined, because according to Bean model $J_c \propto M$ [17]. As can be seen, increase of the film thickness results in enhancement of M (and hence J_c) of the film. The multilayered samples show lower magnetic moment values compared to the YBCO samples at the same thickness and temperature, which is in contradiction to previously studied thicker samples [10,12]. This is believed to occur due to very thin thickness of the individual layers in the multilayered films which likely exhibit an island growth, creating more defects than compared to monolayer YBCO. This leads to a decrease in J_c and magnetic moment.

The electromagnetic homogeneity of the films was investigated by MOI (Fig. 4). From the flux distribution in the Y/Sm44 sample, multiple defects can

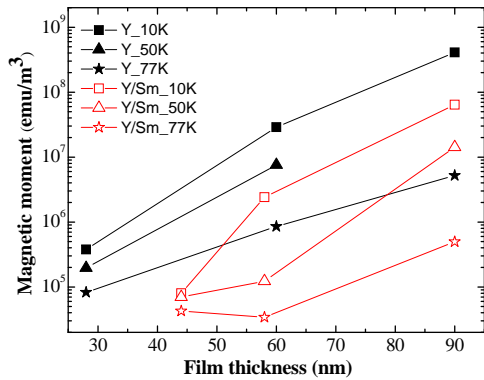


Fig. 3. Magnetic moment of films studied at temperatures of 10 K, 50 K and 77 K. (Lines are guides for eyes only.)

be seen over the entire surface of the film, which allows the magnetic flux (bright areas in pictures) to penetrate the sample more readily. This shows that this sample has poor electromagnetic homogeneity and would explain the low value of magnetic moment obtained from the magnetic measurements.

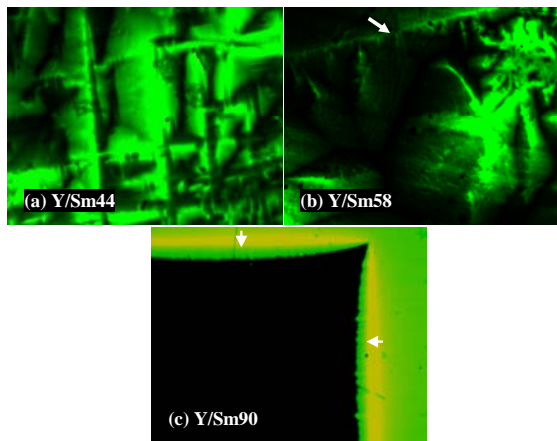


Fig. 4. Magneto-optical images for multilayered YBCO/SmBCO/YBCO samples with different thicknesses. Arrows show the edge of the sample. The external magnetic field is 21 mT for (a) and 33 mT for (b, c).

As the thickness of the film increases (Fig. 4 b, c), dramatically less magnetic flux can penetrate in the sample, showing enhancement of the electromagnetic homogeneity. This reflects the magnetic moment trend, since the lower flux penetration leads to a higher magnetic moment and corresponding J_c .

In summary, the properties of monolayer YBCO and multilayered YBCO/SmBCO/YBCO thin films with thicknesses less than 90 nm were investigated. It was found that as the film's thickness decreased there was a decrease in superconducting character-

istics (T_c , J_c) of the films. The properties obtained were rationalised using SEM and MOI visualisation techniques. We demonstrated that the thinner samples had very poor structural homogeneity, and thus magnetic flux could easily penetrate the sample along many defects. For thicker samples (44 nm - 90 nm), the homogeneity increased dramatically, which is confirmed by enhanced T_c and J_c values and MOI study of the samples. It was also found that multilayering of thin films with thickness ≤ 90 nm results in degradation of their overall superconducting properties. However, there is a qualitative difference in homogeneity between two types of films studied, which suggest better homogeneity of multilayered films taking into account the thickness of each individual layer. We speculate that if epitaxial growth of the initial layers of YBCO and SmBCO is improved, the multilayering approach may enhance structural and electromagnetic properties of the films at very thin (tens of nm) thickness. This may be achieved by utilizing atomic layer engineering or molecular beam epitaxy techniques enabling careful, layer-by-layer control of growing material [14].

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References

- [1] G. N. Gol'tsman et al., Appl. Phys. Lett. 79, 705 (2001);
- [2] H. Shibata et al., Appl. Phys. Lett. 97, 212504 (2010).
- [3] E. Reiger et al., IEEE J. Sel. Top. Quant. Electron. 13, 934 (2007).
- [4] F. Marsili et al., Opt. Expr. 16, 3191 (2008);
- [5] A. Korneev et al., Appl. Phys. Lett. 84, 5338 (2004).
- [6] T. Terashima et al., Phys. Rev. Lett. 65, 2684 (1990).
- [7] T. Terashima et al., Phys. Rev. Lett. 67, 1362 (1991).
- [8] X. Y. Zheng et al., Phys. Rev. B 45, 7584 (1992).
- [9] J. Gao, T. C. Chui, W. H. Tang, IEEE Trans. Appl. Supercond. 9, 1661 (1999).
- [10] A. V. Pan, S. Pysarenko, and S. X. Dou, Appl. Phys. Lett. 88, 232506 (2006).
- [11] O. V. Shcherbakova et al., IEEE Trans. Appl. Supercond. 21, 156 (2011).
- [12] A. V. Pan et al., IEEE Trans. Appl. Supercond. 17, 3585 (2007).
- [13] Q. X. Jia et al., Appl. Phys. Lett. 80, 1601 (2002).
- [14] I. Bozovic, IEEE Trans. Appl. Supercond. 11, 2686 (2001).
- [15] R. T. Thew et al., Nucl. Instrum. Meth. Phys. Res. A 610, 16 (2009).
- [16] N. Curtz et al., Supercond. Sci. Technol. 23, 045015 (2010).
- [17] C. P. Bean, Rev. Mod. Phys. 36, 31 (1964).