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Total electron-scattering cross sections of pyrimidine

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Abstract
Total electron-scattering cross sections of pyrimidine, the basic component for the nucleic bases cytosine and thymine, were measured for electron energies from 5 eV to 1 keV using the linear transmission method. The measured results were compared to semiempirical data obtained by means of the additivity rule and to experimental data for benzene since it has a similar ring structure and the same number of valence electrons as pyrimidine. Furthermore, integral elastic and inelastic electron-scattering cross sections of pyrimidine were calculated by applying the spherical complex optical potential model. The sum of both cross sections agrees reasonably well with the experimental total electron-scattering cross sections of pyrimidine in the energy range from 20 eV to 1 keV. The experimental data are, however, significantly lower than the theoretical cross sections when including the contribution of rotational excitations to the electron scattering.

Keywords
sections, cross, pyrimidine, scattering, total, electron

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Total electron-scattering cross sections of pyrimidine

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It is a general consensus that DNA is the primary target for cellular damage induced by ionizing radiation [1]. When penetrating tissues, ionizing radiation causes cellular damage either by a direct excitation and ionization of the DNA or via secondary electrons that further undergo interactions with the DNA in the lateral environment of the primary radiation track. The secondary electrons are responsible for a significant part of the radiation damage. Their contribution to the radiation damage is predominant in the case of low linear energy transfer (LET) radiations.

Since an experimental investigation of the initial physical action of ionizing radiation at cellular levels is not yet feasible, Monte Carlo simulations are commonly applied to specify the initial radiation damage. For such simulations, the total electron-scattering cross sections (TCS) of DNA constituents play an important role. As the sum of all interaction cross sections, they represent the upper limit for partial scattering cross sections and can therefore be employed to check the consistency of compiled data sets. Total electron-scattering cross sections are also often used to put relative electron-scattering cross sections of molecules which serve as the models of building blocks of DNA into a common scale [2], and in Monte Carlo simulations they are required to determine the mean distance between two subsequent interaction points.

Few experiments have been carried out to determine the TCS of molecules which serve as the models of building blocks of DNA. Up to now, the most attention has been paid to tetrahydrofuran (THF), which is used as the model molecule for the deoxyribose in the backbone of the DNA. Baek et al. [2] determined the TCS of THF for electron energies between 6 eV and 1 keV using the same experimental setup as in the present work. Zecca et al. [3] measured the TCS of THF for low electron energies. Mozejko et al. [4] reported the experimental TCS of THF for electron energies from 1 to 370 eV and Fuss et al. [5] measured those in the intermediate and high energy regions.

Regarding the TCS of nucleic bases for electron scattering, no experimental data have been published yet. Recently, Zecca et al. [6] measured the TCS of pyrimidine, the basic component for the nucleic bases cytosine and thymine, for positrons in the energy range from 0.3 to 45 eV. They calculated the TCS of pyrimidine for electrons in the energy range between 1 eV and 10 keV using the independent-atom screened additivity rule [7,8].

In the present work, the TCS of pyrimidine were measured for electron energies from 5 eV to 1 keV using a linear transmission experiment. The experimental data are compared to theoretical TCS which were calculated by applying the spherical complex optical potential (SCOP) model [9]. Using this model, the integral elastic and inelastic electron-scattering cross sections of pyrimidine were computed for electron energies between 10 eV and 1 keV. Furthermore, the total ionization cross sections of pyrimidine were determined by means of the binary encounter Bethe (BEB) model [10].

II. EXPERIMENT

The measurements were carried out using the linear transmission device described previously [11]. It consists of an electron gun, a scattering chamber, and an electron energy analyzer. The electron gun, equipped with a hairpin tungsten filament, delivered stable electron beams in the energy range from 20 eV to 5 keV. An electrostatic lens and two orthogonal pairs of deflection plates were used to focus and align the electron beam, respectively. The scattering chamber consisted of a hollow metallic cylinder mounted in a vacuum recipient. The symmetry axis of the scattering chamber was set perpendicular to that of the electron beam. Two circular apertures on the cylinder shell defined the electron beam direction. The diameter of both apertures and of the scattering chamber amounted to 0.5 and 132 mm, respectively. The apertures were electrically isolated from the scattering chamber for the monitoring of the change of the incident electron beam current.

To enable the measurement of TCS below 20 eV, the scattering chamber was electrically isolated so that a negative dc bias voltage $U_{ch}$ could be applied [11]. This leads to the reduction of the incident electron energy $T_0$ by the amount of $eU_{ch}$ at the entrance aperture. In this case, the electrons undergo collision processes in the scattering chamber at the reduced energy $T = T_0 - e |U_{ch}|$ and those transmitted...
without interaction regain their initial energy on the path between the exit aperture and the electron energy analyzer. The electron beam underwent collision processes with gaseous pyrimidine in the scattering chamber, leading to an attenuation of the electron current along the beam direction. The electron beam leaving the scattering chamber entered an electron energy analyzer which discriminated unscattered electrons from electrons scattered inelastically in the forward direction. The electron energy analyzer was of hemispherical condenser design with a mean radius of 100 mm and a deflection angle of 150°. A channel electron multiplier, mounted at the end of the analyzer, was used for the detection of electrons.

The geometrical solid angle subtended by the detector varied from $7 \times 10^{-2}$ sr at high electron energies to $1.1 \times 10^{-3}$ sr below 20 eV. The energy resolution (FWHM) was adjusted by changing the retardation voltage applied to the entrance slit of the analyzer [12]. The relative energy resolution was better than 0.25%. The absolute value of the energy resolution was smaller than 0.5 eV at electron energies below 100 eV and amounted to about 1 eV at 1 keV. The whole electron beam path was passively shielded against the residual earth magnetic field as well as other electromagnetic disturbances by means of a permalloy housing. The major part of the external magnetic field was actively compensated by three orthogonal pairs of Helmholtz coils surrounding the scattering chamber.

The scattering chamber was connected to a gas source via a regulating valve with an adjustable leak rate from $10^{-3}$ mbar 1/s to $10^{-1}$ mbar 1/s. Before the introduction of the gas, the scattering chamber was evacuated down to a residual pressure lower than $10^{-6}$ mbar. The vapor of liquid pyrimidine with a stated purity of 99.9% was used as gas source. At room temperatures of higher than 28 °C, the vapor pressure of pyrimidine was sufficiently high for the experiment.

The gas pressure in the scattering chamber was measured by means of a capacitance manometer which had been calibrated at the vacuum metrology section of Physikalisch-Technische Bundesanstalt in Berlin, Germany, in order to correct for the thermal transpiration effect [13]. The calibration was carried out for several light gases such as N$_2$ and CH$_4$ in the pressure range between $10^{-3}$ and 1 mbar. The greatest deviation between the reading of the capacitance manometer used in this work and the primary standard was found in the pressure region around $10^{-3}$ mbar. In the pressure region around 1 mbar, there was almost no difference between both values. Since the typical gas pressure in the present experiment lay between $10^{-3}$ and $10^{-2}$ mbar, a disregard of the thermal transpiration effect would lead to a lowering of the experimental results by about 3%.

As the correction factor $\Gamma$, defined as the ratio of the reading value $p_N$ of the capacitance manometer to the pressure standard $p$, was not known for pyrimidine, it was estimated using the empirical equation for the thermal transpiration effect proposed by Takaishi and Sensui [14]:

$$\frac{1 - (p/p_m)}{1 - (T_1/T_2)^{1/2}} = \frac{1}{(AX^2 + BX + CX^{1/2} + 1)}. \quad (1)$$

$T_1$ and $p$ are the temperature and pressure in the chamber, respectively, and $p_m$ is the pressure in the manometer gauge held at $T_2$ with $T_2 > T_1$. Here, the ratio $p_m/p$ corresponds to the correction factor $\Gamma$. The quantity $X$ is equal to the product of $p_m$ and the diameter $d$ of the connection line between the scattering chamber and the manometer gauge along which a temperature gradient occurs: $X = p_m d$. The constants $A$, $B$, and $C$ are related to gas kinetic collision cross sections. The correction factor for pyrimidine was estimated using the gas kinetic collision cross sections of benzene (C$_6$H$_6$) whose molecular diameter and structure is similar to that of pyrimidine. Here, the gas kinetic collision cross section of benzene was determined from its dynamic viscosity value [15]. Figure 1 shows the estimated correction factor for pyrimidine in comparison to that of methane.

The determination of $\text{TCS} \sigma_t$, was based on Beer’s attenuation law taking into account the secondary effects:

$$c = c_0 \left[ \exp(-\sigma_t n_F) - \int n(z) dz \right] = \int \left[ \frac{d\sigma}{d\Omega} \right] dz d\Omega, \quad (2)$$

where $c_0$ and $c$ are the electron count rate of the channel electron multiplier without and with the gas in the scattering chamber, respectively. The quantity $n_F$ is the number of molecules per area crossed by the electron beam and is given, when considering the inhomogeneity of the number density due to the gas streaming out through the apertures, by the integral of the number density distribution $n(z)$ over the electron beam trajectory:

$$n_F = \int n(z) dz \equiv k_1 n_0 L_0. \quad (3)$$

In Eq. (3), $n_0$ is the number density in the static case where there is no gas flow through the apertures and $L_0$ is the diameter of the scattering chamber. In the static case, $n_0$ can be determined from the ideal gas law $n_0 = p/(k_B T_1)$, where the gas temperature $T_1$ is assumed to be equal to the room temperature measured with a uncertainty smaller than 1 K.

The second integral on the right-hand side of Eq. (2) corrects for the signal contribution from electrons that underwent scattering processes but were counted due to the finite

FIG. 1. Pressure correction factor $\Gamma$ for pyrimidine as a function of the reading $p_N$ of the capacitance manometer. For comparison, the experimental corrections factors ($\circ$) of methane are shown.
The elastic-scattering cross section $d\sigma/d\Omega$ depends on the scattering angle and on the electron energy. In the present work, $d\sigma/d\Omega$ was set to be equal to the differential elastic-scattering cross section $d\sigma_{el}/d\Omega$; in other words, the rotational and vibrational excitations whose transition energies are smaller than the energy resolution of the electron energy analyzer were not taken into account in $d\sigma/d\Omega$. The consequence of disregarding these excitation processes is discussed below in the comparison of the experimental data with theoretical values. The $z$-dependent detector solid angle $\Delta\Omega(z)$ can be approximated by an average value, as the distance between the center of the scattering chamber and the detector, which was 384 mm, was much greater than the diameter of the scattering chamber:

$$\int n(z)dz \int _{\Delta\Omega(z)} \frac{d\sigma_{el}}{d\Omega} \approx \frac{d\sigma_{el}}{d\Omega} \Delta\Omega \int n(z)dz$$

$$= \frac{d\sigma_{el}}{d\Omega} \Delta\Omega k_1 n_0 L_0 = \Delta\sigma_{el} k_1 n_0 L_0. \quad (4)$$

In Eq. (4), $\Delta\Omega$ is the detector solid angle subtended at the position $z = L_0/2$. In the present experimental arrangement, $\Delta\sigma_{el}$ is in the order of 1% of $\sigma$. Therefore, Eq. (2) can be approximately rewritten as

$$\ln \left( \frac{c}{c_0} \right) \approx k_1 n_0 L_0 (\sigma_0 - \Delta\sigma_{el}). \quad (5)$$

According to Eq. (5), the TCS $\sigma_t$ can be determined by recording the count rate $c$ as a function of the gas number density $n_0$ or the chamber pressure $p$, if the values of $k_1$ and $\Delta\sigma_{el}$ are known. The correction factor $k_1$, which accounts for the pressure drop in the vicinity of the apertures and for the additional beam attenuation in the effusing gas stream, was estimated using the calculations of Nelson and Colgate [16]. The resulting value of $k_1$ for the apertures and the scattering chamber used in this work amounted to 0.99.

The value of $\Delta\sigma_{el}$ was estimated by means of the differential elastic electron-scattering cross sections $d\sigma_{el}/d\Omega$ calculated using the modified independent-atom model [17]. In the present work, the pressure in the scattering chamber varied between the residual pressure of about $1.0 \times 10^{-6}$ mbar and a maximal pressure depending on the electron energy. The maximal pressure value was adjusted such that the beam chamber used in this work amounted to 0.99.

A considerable increase of the electron beam current was observed when pyrimidine vapor was introduced into the scattering chamber. This increase arises due to the chemical reaction of the tungsten filament with the pyrimidine gas streaming out through the entrance aperture of the scattering chamber. In order to take into account the change of the electron emission due to the gas exposure of the filament, the primary electron beam current $I_p$ was measured on the entrance aperture of the scattering chamber. The TCS $\sigma_t(T)$ was then obtained by a linear fit of $\ln([c/I_p]/[c_0/I_{p,0}])$ versus $n_0 L_0$, where $I_{p,0}$ is the primary beam current if no gas is present in the scattering chamber. The slope $\sigma_0$ of the linear fit is given by Eq. (5) and equal to $k_1(\sigma_0 - \Delta\sigma_{el})$.

The measurement was repeated ten times for each electron energy to reduce the uncertainties arising from statistical fluctuations. The TCS $\sigma_t$ of pyrimidine was then determined using the mean value $\bar{\sigma}_0$ of the ten measurements:

$$\sigma_t = \sigma_0/k_1 + \Delta\sigma_{el}. \quad (6)$$

The standard uncertainty $u(\sigma_t)$ of the experimental TCS was calculated according to the Guide to the Expression of Uncertainty in Measurements [18]:

$$u^2(\sigma_t) = \left( \frac{\partial\sigma_t}{\partial k_1} \right)^2 + \left( \frac{\partial\sigma_t}{\partial \bar{\sigma}_0} \right)^2 + \left( \frac{\partial\sigma_t}{\partial (\Delta\sigma_{el})} \right)^2 \left( \Delta(\Delta\sigma_{el}) \right)^2 + \left( \frac{\partial\sigma_t}{\partial T} \right)^2 \delta T^2, \quad (7)$$

where only the main uncertainty sources are taken into account. The uncertainty $\Delta k_1$ amounted to $1 \times 10^{-2}$ and the relative value of the uncertainty $\delta\sigma_0$ varied between 1% and 10% depending on the electron energy. The uncertainty of $\Delta\sigma_{el}$ is mainly given by that of $d\sigma_{el}/d\Omega$ of pyrimidine. The relative uncertainty $\delta(\Delta\sigma_{el})/\Delta\sigma_{el}$, estimated by comparing the calculated $d\sigma_{el}/d\Omega$ with available experimental data [19,20], amounted to 50%. The derivative $\partial\sigma_t/\partial T$ was determined using the cubic interpolation of the experimental results and $\delta T$ was set equal to the uncertainty of the calibration of the electron energy that has been carried out before by means of the retarding field method [11,21].

### III. THEORETICAL METHODS

In addition to the measurement, the TCS of pyrimidine were calculated by means of the SCOP model [9]. This model is based on the fixed-nuclei approximation and describes the electron scattering by means of a complex potential $V_{opt}(\vec{r})$ given by

$$V_{opt}(\vec{r}) = V_R(\vec{r}) + i V_{abs}(\vec{r}), \quad (8)$$

where the real and imaginary parts of the potential account for elastic and inelastic scattering, respectively. The real part $V_R(\vec{r})$ consists of the static potential $V_{el}(\vec{r})$, the exchange potential $V_{ex}(\vec{r})$, and the correlation-polarization potential $V_{pol}(\vec{r})$:

$$V_R(\vec{r}) = V_{el}(\vec{r}) + V_{ex}(\vec{r}) + V_{pol}(\vec{r}). \quad (9)$$

The imaginary part $V_{abs}(\vec{r})$ describes the dissipation of incoming electron waves into the inelastic channels such as ionization and excitation.

In the SCOP model, $V_{opt}(\vec{r})$ is at first expanded around the center of mass of the molecule using symmetry-adapted functions [22]:

$$V_{opt}^{\mu}(\vec{r}) = \sum_{l,h} \psi_{lh}(\hat{r}) X_{lh}^{\mu}(\vec{r}), \quad (10)$$

where $\mu$ designates the particular irreducible representation of the molecular point group in the ground state, $h$ denotes a specific basis at an angular momentum $l$ for $p\mu$, $r$ is the distance from the center of mass of the molecule, $\hat{r}$ is the unit vector of $r$, and $\psi_{lh}$ is the corresponding expansion coefficient. For closed-shell nonlinear molecules such as pyrimidine, the ground-state symmetry $p\mu$ is equal to the irreducible...
representation $^1 A_1$ and the symmetry-adapted function $X_{1h}^{A_1}(\hat{r})$ can be represented as a linear combination of real spherical harmonics $S_{lm}(\hat{r})$ [23]:

$$X_{1h}^{A_1}(\hat{r}) = \sum_{m=0}^{l} b_{lm}^{A_1} S_{lm}(\hat{r}).$$  \hspace{1cm} (11)

The values of the coefficients $b_{lm}^{A_1}$ are given by the character table of the corresponding irreducible representation.

The main feature of the SCOP model [9] is that only the isotropic term in the multipole expansion, i.e., the first term with $l = 0$ and $h = 0$, is considered. In this case, the optical potential $V_{\text{opt}}^{\text{SCOP}}(\hat{r})$ can be written as

$$V_{\text{opt}}^{\text{SCOP}}(\hat{r}) = v_{01}(r)b_{010}^{A_1} S_{00}(\hat{r}) = v_{01}(r)/\sqrt{4\pi}. \hspace{1cm} (12)$$

As the following equation shows, the consideration of only the isotropic term is equivalent to averaging $V_{\text{opt}}(\hat{r})$ over the entire solid angle:

$$\bar{V}_{\text{opt}}(r) = \frac{\int V_{\text{opt}}(\hat{r})d\Omega}{4\pi} = \frac{\sqrt{4\pi} \int V_{\text{opt}}(\hat{r})S_{00}(\hat{r})d\Omega}{4\pi} = \frac{\int \sum_{l,h} v_{01}(r) \sum_{m=0}^{l} b_{lm}^{A_1} S_{lm}(\hat{r}) S_{00}(\hat{r})d\Omega}{\sqrt{4\pi}} = \frac{v_{01}(r)}{\sqrt{4\pi}} = V_{\text{opt}}^{\text{SCOP}}(r), \hspace{1cm} (13)$$

where use was made of the orthogonality relation of the spherical harmonics $\int S_{lm}(\hat{r}) S_{lm'}(\hat{r})d\Omega = \delta_{ll'}\delta_{mm'}$. The above averaging can be interpreted as the averaging of the interaction potential over all molecular orientations which are randomly distributed in the gas beam.

All four potentials given on the right-hand side of Eqs. (8) and (9) depend on the electron density distribution $\rho(\hat{r})$, which was calculated by means of the quantum chemical program GAUSSIAN 09 [24] deploying the 6-311G* basis sets. The static and correlation-polarization potentials $V_{\text{el}}(\hat{r})$ and $V_{\text{pol}}(\hat{r})$, respectively, were then computed from this electron density distribution with the help of the library SCELIB [23]. It should be noted that the calculation of $V_{\text{pol}}(\hat{r})$ was based on the formulations for the VCP polarization potential given in Ref. [23]. The exchange potential $V_{\text{ex}}(\hat{r})$ was obtained by the application of the following formula which was derived from the free electron gas model of Hara [25]:

$$V_{\text{ex}}(\hat{r}) = -\frac{2}{\pi} k_F(\hat{r}) \left( \frac{1}{2} + \frac{1 - \xi^2}{4\xi} \ln \left| 1 + \frac{1 - \xi}{1 - \zeta} \right| \right) \hspace{1cm} (14)$$

with $k_F(\hat{r}) = [3\pi^2 (\rho(\hat{r}))]^{1/3}$ and $\zeta(\hat{r}) = (k^2 + 2I + k_F^2)^{1/2}/k_F$, where $k$ is the initial momentum of the incident electron and $I$ is the ionization potential in atomic units.

The absorption potential that is dependent on the electron energy $T$ was determined using the quasifree-scattering model described in detail by Staszeweska et al. [26,27]. In this model, $V_{\text{abs}}(\hat{r})$ is given by

$$V_{\text{abs}}(\hat{r}) = -\rho(\hat{r})\sqrt{2(T - V_{\text{ex}})} \left( \frac{4\pi}{10k_F^2 T} \right) H \times (k^2 - k_F^2 - 2\Delta)(A_1 + A_2 + A_3), \hspace{1cm} (15)$$

where $V_{\text{ex}}(\hat{r}) = V_{\text{el}}(\hat{r}) + V_{\text{ex}}(\hat{r})$, $H(x)$ is the Heaviside step function, such that $H(x) = 1$ for $x \geq 0$ and $H(x) = 0$ for $x < 0$, and $\Delta$ is the mean excitation energy. The terms $A_1$, $A_2$, and $A_3$ are defined as

$$A_1 = \frac{5k_F^3}{2\Delta}, \hspace{1cm} A_2 = \frac{k_F^3 (5k_F^2 - 2k_F^2)}{(k^2 - k_F^2)} \hspace{1cm} (16)$$

$A_3 = H(2k_F^2 + 2\Delta - k^2) \left( 2(k_F^2 + 2\Delta - k^2)^{3/2} \right) / (k^2 - k_F^2)$.

Jain and Baluja [9] calculated the mean excitation energy $\Delta$ for a variety of molecules and found that the value of $\Delta$ is very close to the ionization potential in the most cases. Therefore, $\Delta$ was set equal to the ionization potential $I$ of pyrimidine.

Figure 2 depicts the spherically averaged potentials $\bar{V}_{\text{el}}(r)$, $\bar{V}_{\text{ex}}(r)$, $\bar{V}_{\text{pol}}(r)$, and $\bar{V}_{\text{abs}}(r)$ as functions of the distance from the center of mass of the molecule.

As $V_{\text{opt}}(r)$ is spherically symmetric, the wave function of the scattered electrons can be determined by solving radial Schrödinger equations in combination with the partial wave expansion method:

$$\left[ \frac{d^2}{dr^2} + k^2 - \frac{l(l+1)}{r^2} - 2\bar{V}_{\text{opt}}(r) \right] \psi_l(r) = 0, \hspace{1cm} (17)$$

where $k$ is the momentum of the incident electron and $\psi_l$ is the partial wave for the angular momentum quantum number $l$.

Equation (17) was solved by applying the variable phase approach [28], which transforms the complex second-order differential equation into two coupled nonlinear first-order differential equations. In this approach, the real part $\epsilon_l(kr)$ and the imaginary part $\eta_l(kr)$ of the phase shift can be obtained from the following system of coupled differential equations:

$$\epsilon'_l = -\frac{1}{k} \left[ \bar{V}_R(X^2 - Y^2) - 2\bar{V}_{\text{abs}}XY \right], \hspace{1cm} \eta'_l = -\frac{1}{k} \left[ \bar{V}_{\text{abs}}(X^2 - Y^2) + 2\bar{V}_R XY \right] \hspace{1cm} (18)$$
with
\[ X = \cosh \eta (\hat{n}_1 \sin \epsilon_1 - \hat{j}_1 \cos \epsilon_1), \]
\[ Y = \sinh \eta (\hat{n}_1 \cos \epsilon_1 + \hat{j}_1 \sin \epsilon_1), \]  \tag{19}
where the phase shift fulfills the boundary conditions \( \epsilon_1(0) = 0 \) and \( \eta_1(0) = 0 \). The functions \( \hat{n}_1(kr) \) and \( \hat{n}_1(kr) \) are the Riccati-Bessel functions. It should be noted that Eqs. (18) and (19) deviate from those found in literature [9, 22, 29–32] which, in contrast to our Eqs. (18) and (19), are inconsistent with the original formulas given in Ref. [28].

The phase shifts \( \epsilon_1(kr) \) and \( \eta_1(kr) \) for \( r \to \infty \) are related to the \( S \) matrix \( S_l(k) \) via
\[ S_l(k) = \exp[-2\eta_l] \exp[2i\epsilon_l]. \]  \tag{20}

The integral elastic \( \sigma_{\text{el}} \), absorption \( \sigma_{\text{abs}} \), and total \( \sigma_j \) cross sections can be calculated from the \( S \) matrix \( S_l(k) \) by means of the following equations:
\[ \sigma_{\text{el}}(k) = \frac{\pi}{k^2} \sum_{l=0}^{l_{\text{max}}} (2l + 1) |1 - S_l(k)|^2, \]
\[ \sigma_{\text{abs}}(k) = \frac{\pi}{k^2} \sum_{l=0}^{l_{\text{max}}} (2l + 1) [1 - |S_l(k)|^2], \]
\[ \sigma_j(k) = \frac{2\pi}{k^2} \sum_{l=0}^{l_{\text{max}}} (2l + 1) [1 - \text{Re} S_l(k)]. \]  \tag{21}

The values of \( S_l(k) \) were obtained from the asymptotic values of \( \epsilon_l(kr) \) and \( \eta_l(kr) \) via Eq. (20), where \( \epsilon_l(k) \) and \( \eta_l(k) \) were determined by the numerical integration of Eq. (18) over \( r \) using the fourth-order Runge-Kutta method [33]. The integration was carried out until the relative change of the absolute value of the phase shift was smaller than \( 1.0 \times 10^{-5} \).

As mentioned above, the SCOP model [9] takes into account only the isotropic \( l = 0 \) term in the multipole expansion of \( V_{\text{opt}}(r) \). The disregard of the anisotropic terms with \( l \geq 1 \), however, can lead to a considerable underestimate of electron-scattering cross sections in the case of molecules with high dipole moments like pyrimidine since such molecules can be rotationally excited during the interaction with incident electrons.

Following the approach used by Jain [34], the contribution of rotational excitations to the electron-scattering cross section was calculated using the formula of Collins and Norcross [35] and incoherently added to the integral cross sections obtained using Eq. (21). According to the formula of Collins and Norcross [35], which is based on the first Born approximation for a rotating dipole, the integral cross section for rotational excitation is given by
\[ \sigma_{\text{rot}}(j \to j') = \frac{8\pi}{3k^2} D_m^2 \frac{j'_1}{2j + 1} \ln \frac{k + k'}{|k - k'|}, \]  \tag{22}
with
\[ 2k^2 = 2k^2 - (E_{j'\tau} - E_{j\tau}), \]  \tag{23}
where \( D_m \) is the molecular dipole moment and \( E_{j\tau} \) is the energy of the rotational state \( (j \tau) \). The energy levels of the rotational states of pyrimidine, which is an asymmetric top molecule, were calculated by means of the computational method described in Ref. [36]. Assuming that the dipole transition to the first rotational state \( (00 \to 00) \), whose excitation energy amounts to \( E_{10} - E_{00} = 4.0 \times 10^{-2} \) meV, is the major transition, the electron-scattering cross section \( \sigma_{\text{rot}} \) due to rotational excitation was calculated using Eqs. (22) and (23) and added to \( \sigma_j \) calculated using Eq. (21).

### IV. RESULTS AND DISCUSSION

The results of the present measurements are shown in Fig. 3 and listed in Table I. Apart from the weak shoulderlike

<table>
<thead>
<tr>
<th>( T ) (eV)</th>
<th>( \sigma_j ) ( (10^{-16} \text{ cm}^2) )</th>
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<tr>
<td>5</td>
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<tr>
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<td>8</td>
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<td>200</td>
<td>18.73 ± 0.68</td>
</tr>
<tr>
<td>300</td>
<td>14.82 ± 0.98</td>
</tr>
<tr>
<td>400</td>
<td>12.76 ± 0.43</td>
</tr>
<tr>
<td>600</td>
<td>9.78 ± 0.63</td>
</tr>
<tr>
<td>800</td>
<td>8.13 ± 0.70</td>
</tr>
<tr>
<td>1000</td>
<td>6.94 ± 0.39</td>
</tr>
</tbody>
</table>

TABLE I. Total electron-scattering cross section \( \sigma_j \) of pyrimidine.
The additivity rule \[41\]: (comparison to the semiempirical values (Fig. 4) compared to the semiempirical values which TCS in the low-energy region.

greater than that of benzene, leads to a strong elevation of the dipole moment of pyrimidine, which is one order of magnitude low electron energies is more realistic as the relatively high should be noted that the disagreement between both data at the TCS of pyrimidine at electron energies below 50 eV. It should be noted that the disagreement between both data at low electron energies is more realistic as the relatively high dipole moment of pyrimidine, which is one order of magnitude greater than that of benzene, leads to a strong elevation of the TCS in the low-energy region.

In addition to the TCS of benzene, the present results are in Fig. 4 compared to the semiempirical values which were obtained using the TCS of C\(_2\)H\(_2\) [39], N\(_2\) [40], and the additivity rule [41]. As can be seen from Fig. 4, the experimental and semiempirical values agree within the experimental uncertainties over the whole energy range, even at electron energies lower than 50 eV. It should be noted that the disagreement between both data at low electron energies is more realistic as the relatively high dipole moment of pyrimidine, which is one order of magnitude greater than that of benzene, leads to a strong elevation of the TCS in the low-energy region.

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Apart from the shoulderlike structure of the elastic-scattering cross section at electron energies around 10 eV, the TCS of pyrimidine decreases with increasing electron energy. Since there are no experimental data for the TCS of pyrimidine measured by other groups, the present results are compared to the TCS of benzene. The TCS of both molecules are expected to agree approximately at the intermediate and high electron energies because they are isoelectronic and have a similar structure. It is evident from Fig. 3 that the present results are well reproduced by the benzene data of Możejko et al. [37], even at low electron energies, while those of Sueoka [38] are noticeably lower than the TCS of pyrimidine at electron energies below 50 eV. It should be noted that the disagreement between both data at low electron energies is more realistic as the relatively high dipole moment of pyrimidine, which is one order of magnitude greater than that of benzene, leads to a strong elevation of the TCS in the low-energy region.

The values of the binding energy \(B_i\) and of the average electron kinetic energy \(U_i\) were determined by means of the quantum chemical code GAUSSIAN 09 [24] deploying the 6-311G** set.

The inelastic scattering and ionization cross sections rise up to a maximal value located around 80 eV and then decrease monotonically with increasing energy. As can be seen from Fig. 5, the difference between the inelastic scattering and total ionization cross section, corresponding to the excitation cross section, decreases with increasing electron energy, in accordance with the fact that inelastic scattering at high electron energies occurs predominantly due to ionization processes.

Figure 6(a) shows the present experimental results in comparison to the theoretical TCS of pyrimidine obtained by summing the elastic- and inelastic-scattering cross sections depicted in Fig. 5. The solid and dashed lines represent the theoretical values with and without the contribution of rotational excitations to the electron scattering, respectively.
additivity rule [7,8] and also the Born approximation in high electron energies. It is also evident from Fig. 6(a) between both data decreases with increasing energy to vanish than the experimental results. Above 30 eV, the deviation 30 eV where the theoretical values are about 50% higher and theoretical data was found at electron energies around 10 eV. The greatest deviation between the present experimental by the results of Zecca

The good agreement of the theoretical TCS of benzene with the measured results of the Gdański group [37] whose experimental setup was similar to that of this work supports the assumption that the discrepancy between the present theoretical and experimental data for pyrimidine may mainly arise due to insufficient resolving power of the present apparatus with respect to the electron-scattering processes caused by rotational excitations. Indeed, the angular distribution of these processes is sharply peaked in the forward direction. According to the estimate carried out using the formula for the differential cross section given in Ref. [35], the part of electrons scattered due to rotational excitations within the detection solid angle of the present experiment makes up about 60%–70% of the total rotational excitation cross section at electron energies around 30 eV. Apart from rotational excitations, nonresolved vibrational transitions and the counting of electrons scattered at metallic surfaces such as slits and apertures which can result in an effective detection solid angle being larger than the geometrically defined one can lead to a decrease of the measured TCS.

V. CONCLUSION

Total electron-scattering cross sections of pyrimidine were measured for electron energies from 5 eV to 1 keV by means of a linear transmission device. The results of this work agree surprisingly well with the semiempirical values obtained using the additivity rule [41] and with the experimental TCS of the isoelectronic molecule benzene determined by Mozejko et al. [37] over the whole measured energy range. The good agreement, particularly at low electron energies, is remarkable since one expects the relatively high permanent dipole moment of pyrimidine to strongly increase its TCS with decreasing electron energy when compared to the case of the TCS of benzene and of the semiempirical values.

The present experimental data are rather well reproduced by the theoretical values of Zecca et al. [6] that were obtained.
using the independent-atom screened additivity rule \cite{7,8} including the contribution of rotational excitations to the electron scattering. On the other hand, the theoretical TCS of pyrimidine calculated with the SCOP model \cite{9} that employs electron scattering. On the other hand, the theoretical TCS including the contribution of rotational excitations to the theoretical results in the case of benzene suggest that the electrons scattered due to rotational excitations may be incompletely filtered out in the present measurement. For the precise measurement of the TCS of molecules with high permanent dipole moments, an experimental setup with a high angular and energy resolving power or an accurate determination of \( \Delta \sigma_{\alpha} \) including rotational excitations is therefore envisaged.

**ACKNOWLEDGMENT**

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\[ \text{dipole moments, an experimental setup with a high angular} \]
\[ \text{and energy resolving power or an accurate determination of} \]
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\[ \text{envisaged.} \]